

1 **Transport structure of the South Atlantic Ocean derived from a high-resolution**
2 **numerical model and observations**

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16 **Key Points:**

17 1. The upper limb of the AMOC originates from the warm Indian Ocean water through the

18 Agulhas Leakage

19 2. The NADW in the lower limb of AMOC flows southward as a DWBC to 45S and then turns

20 eastward to flow across the Mid-Atlantic Ridge near 42S

21 3. The AMOC variability at seasonal to decadal timescales is latitudinally coherent from 35S

22 to about 35N

23 **Abstract**

24 The South Atlantic Ocean plays an important role in the Atlantic meridional overturning
25 circulation (AMOC), connecting it to the Indian and Pacific Oceans as part of the global
26 overturning circulation system; yet, the detailed time mean circulation structure in this
27 region and the large-scale spatial pattern of the AMOC variability remain unclear. Using
28 model outputs from a 50-year, eddying global ocean-sea ice simulation validated against
29 observations at a zonal section at 34°S, a meridional section at 65°W in the Drake Passage,
30 and a meridional section southwest of Africa, we find that the upper limb of the AMOC
31 originates from the Agulhas leakage and that while the cold Pacific water from the Drake
32 Passage does not contribute directly to the AMOC it does play a significant role in setting the
33 temperature and salinity properties of the water in the subtropical South Atlantic. We also
34 find that the North Atlantic deep water (NADW) in the lower limb of AMOC flows southward
35 as a deep western boundary current all the way to 45°S and then turns eastward to flow
36 across the Mid-Atlantic Ridge near 42°S, and that the recirculation around the Vitoria-
37 Trindade seamount chain brings some NADW into the Brazil Basin interior. Finally, we find
38 that the AMOC variability is coherent on seasonal to decadal timescales from 35°S to about
39 35°N, where diapycnal water mass transformations between the upper and lower limbs of
40 the AMOC are expected to be small.

41 **1. Introduction**

42 In the Atlantic Ocean, warm water from the South Atlantic flows northward in approximately
43 the upper 1000 meters, loses buoyancy to the atmosphere by cooling *en route* to the northern
44 North Atlantic, and eventually sinks and returns southward at depth as the cold North
45 Atlantic Deep Water (NADW). The temperature difference between the upper and lower
46 limbs of this Atlantic meridional overturning circulation (AMOC) leads to a large northward
47 oceanic heat transport throughout the entire Atlantic basin, in contrast to the poleward heat
48 transport in the Indo-Pacific Ocean (e.g., Macdonald and Baringer, 2013). The South Atlantic
49 Ocean, defined here as the area south of 20°S (Figure 1), plays an important role in that it is
50 through this region where the upper and lower AMOC limbs are connected to the Indian and
51 Pacific Oceans and are entangled in the global overturning circulation system (e.g., Gordon,
52 1986; Broecker, 1991; Schmitz, 1995, 1996; Richardson 2008; Talley, 2013). Thus, a
53 comprehensive knowledge of the circulation in this region is essential to our understanding
54 of the spatial structure and temporal variability of the AMOC.

55 Significant observations have been made in the last 10 years or so toward quantifying and
56 monitoring the AMOC in the South Atlantic, particularly along a latitude near 34.5°S (e.g.,
57 Baringer and Garzoli, 2007; Dong et al., 2009, 2014, 2015; Garzoli et al., 2013; Goes et al.,
58 2015; Meinen et al., 2013; 2018). These observations, which consist of moorings, expendable
59 Bathythermograph (XBT), and Argo float measurements, yield a time mean AMOC transport
60 on the order of 14-20 Sv. They also show that there is significant AMOC variability on several
61 timescales, similar to that observed by the RAPID array at 26.5°N (e.g., McCathy et al., 2015).
62 Beyond 34.5°S, however, the observations in the South Atlantic remain sparse and short (in

63 time). Overall, our understanding of the spatial structure of the time mean circulation is
64 mostly limited to the schematic of Stramma and England (1999) and even less is known
65 about its temporal variability.

66 In particular, there is a long-standing debate regarding the source of the upper limb of the
67 AMOC (Gordon, 2001); whether it originates from the warm, saline Indian waters through
68 the southern rim of Africa (e.g., Gordon 1986; Saunders and King 1995) or from the cooler,
69 fresher Pacific water through the Drake Passage (e.g., Rintoul, 1991; Schlitzer, 1996).
70 Although recent studies seem to favor the warm water route from the Indian Ocean through
71 the Agulhas leakage (e.g., Richardson 2007; Beal et al., 2011), the relative contributions of
72 cold versus warm water is still uncertain (Garzoli and Matano, 2011). At depth, in the lower
73 limb of the AMOC, much attention has been paid to an eastward flow of the NADW near 22°S
74 (e.g., Speer et al., 1995; Stramma and England, 1999; Arhan et al., 2003; Hogg and Thurnherr,
75 2005; Garzoli et al., 2015). This eastward transport is relatively minor when compared to
76 the main branch of NADW, which continues to flow southward as the deep western boundary
77 current (DWBC) along the deep continental slope of South America. However, the exact
78 location of where the DWBC turns eastward south of 34°S and flows across the Mid-Atlantic
79 Ridge (MAR) is largely unknown. Finally, there is also the underlying question as to whether
80 the AMOC variability is meridionally coherent throughout the whole Atlantic (Kelly et al.,
81 2004; Xu et al., 2014).

82 Three-dimensional circulation information beyond the existing observations is required in
83 order to address the above questions. In this paper, we use a high-resolution numerical
84 model to investigate the structure of the mean circulation in the South Atlantic as well as the
85 temporal variability of the AMOC on basin scales. The paper is structured as follows: Section

86 2 summarizes the basic features of the numerical simulation. Section 3 examines the
87 modeled large-scale circulation pattern and the circulation structure at three key
88 observation locations: 34°S across the South Atlantic, 65°W in the Drake Passage, and along
89 a Prime Meridian-Good Hope section southwest of Africa (Figure 1). The model results are
90 shown to be in good agreement with the observed transports and are then used to document
91 the time mean circulation pattern in the South Atlantic (Section 4) and the latitudinal
92 coherence of the AMOC variability throughout the Atlantic basin (Section 5). Summary and
93 discussions follow in Section 6.

94 **2. Numerical Simulation**

95 The numerical results presented in this study are from a long-term global ocean-sea ice
96 hindcast simulation performed using the Hybrid Coordinate Ocean Model (HYCOM, Bleck,
97 2002; Chassignet et al., 2003), coupled with the Community Ice Code (CICE, Hunke and
98 Lipscomb, 2008). The vertical coordinate of the HYCOM is isopycnic in the stratified open
99 ocean and makes a dynamically smooth and time-dependent transition to terrain following
100 in the shallow coastal regions and to fixed pressure levels in the surface mixed layer and/or
101 unstratified seas. In doing so, the model combines the advantages of the different coordinate
102 types in simulating coastal and open ocean circulation features simultaneously (e.g.,
103 Chassignet et al., 2006).

104 The simulation has a horizontal resolution of $1/12^\circ$ (~ 6 km in the area of interest) and a
105 vertical resolution of 36 layers (in σ_2). It is initialized using the January temperature and
106 salinity from an ocean climatology (Carnes, 2009) and is forced using the 3-hourly, 0.7°
107 DRAKKAR forcing set DFS5.2 (Dussin et al., 2016). The DFS5.2 combines two ECMWF

108 (European Centre for Medium-range Weather Forecast) reanalysis products, ERA40 and
 109 ERA-interim, and covers the time period of 1958-2015. The surface heat flux forcing is
 110 computed using the shortwave and longwave radiations from DFS5.2, as well as the latent
 111 and sensible heat fluxes derived from the CORE bulk formulae of Large and Yeager (2004)
 112 and the model sea surface temperature (SST). The surface freshwater forcing includes
 113 evaporation, precipitation, and climatological river runoffs. In addition, the model sea
 114 surface salinity (SSS) is restored toward ocean climatology with a restoring timescale of two
 115 months and it is constrained by an ad hoc assumption of zero global net flux at each time
 116 step. The wind stress is calculated from the atmospheric wind velocity and does not take into
 117 account the shear introduced by the ocean currents. The simulation starts from rest and is
 118 integrated over 1958-2015 with no data assimilation. The horizontal diffusion parameters
 119 are listed in Table 1. In this study, we focus on the last 35 years of the simulation (1981-
 120 2015) as being representative of the time-mean circulation after spin-up (see Figure 2
 121 below).

122

123 Table 1. Viscosity and diffusion coefficients used in the 1/12 global ocean simulation

Parameters	Values
Laplacian coefficient for momentum	$20 \text{ m}^2 \text{ s}^{-1}$
Biharmonic diffusive velocity for momentum	2 cm s^{-1}
Biharmonic diffusive velocity for layer thickness	2 cm s^{-1}
Laplacian diffusive velocity for tracers	0.5 cm s^{-1}

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125 Although a detailed evaluation of the global ocean circulation and the sea ice is beyond the
126 scope of this work, some basic measures are useful. Figure 2 displays the evolution of the
127 domain averaged potential temperature and total kinetic energy of the world ocean. The
128 model temperature exhibits a warming of $\sim 0.02^\circ\text{C}$ per decade (Figure 2a, domain-averaged
129 salinity is constant because of the zero net freshwater flux constraint). The warming is
130 induced by a positive net heat flux of $\sim 1.2\text{ W/m}^2$ (averaged over 1958-2015). This net heat
131 flux in model is about twice the estimates of $0.50 (\pm 0.43)\text{ W/m}^2$ based on the energy gain at
132 the top of atmosphere for 2001-2010 (Loeb et al., 2012) and $0.64 (\pm 0.11)\text{ W/m}^2$ based on 0-
133 700 m oceanic heat content change from combined XBT and Argo data for 1993-2008
134 (Roemmich et al., 2015). However, the modeled net heat flux is within the range of 16 air-
135 sea heat flux products based on global ocean/coupled reanalysis, which found an ensemble
136 mean net heat flux of 1-2 W/m^2 in 1993-2009 ($4.2 \pm 1.1\text{ W/m}^2$ without data-assimilation)
137 (Valdivieso et al., 2017). The modeled total kinetic energy spins up quickly to $\sim 35\text{ cm}^2/\text{s}^2$ in
138 the first two years and remains on that level for the rest of the integration (Figure 2b). The
139 global $1/10^\circ$ simulation performed by Maltrud and McClean (2005) using the Parallel Ocean
140 Program (POP) also reached a maximum in kinetic energy in the first couple of years, except
141 that it is weaker and levels off at about $25\text{-}30\text{ cm}^2/\text{s}^2$ (their Figure 1).

142 The sea ice is quantified and monitored in term of sea ice extent, defined as the area with
143 15% or higher sea ice concentration. Figure 3 compares the evolution of the modeled sea ice
144 extent in million km^2 with the updated results from the National Snow and Ice Data Center
145 (Fetterer et al., 2017). There is a general agreement between model and data in both the
146 northern and southern hemispheres and on both seasonal (Figure 3a) and interannual
147 (Figure 3b) timescales. In particular, as in the observations, the modeled sea ice extent has

148 been decreasing in the northern hemisphere since the beginning of the observations in 1979
149 (note that the pace of the decline is slightly slower in the model). The modeled sea ice extent
150 in the southern hemisphere is relatively stable or increases slightly over time and is in good
151 agreement with the observations (Figure 3b).

152 Figure 4 shows the modeled AMOC transport at 26.5°N , defined as the northward trans-basin
153 transport above the modeled time mean maximum overturning depth (1000 m). The
154 modeled AMOC transport starts at 15 Sv and exhibits a decreasing trend of 0.43 Sv per
155 decade (Figure 4a) with a mean value of 12.2 Sv in 2004-2015. The latter value is 4.6 Sv or
156 ~30% lower than the 16.8 Sv based on the RAPID data (e.g., McCarthy et al., 2015). A similar
157 low AMOC transport of 12.4 Sv is also found in a $1/12^{\circ}$ global ocean simulation using the
158 same atmospheric forcing over the same time period but with a different ocean model, i.e.
159 NEMO (Blaker et al., 2015). The low AMOC may therefore be due to our choice of atmospheric
160 forcing. The modeled AMOC variability, however, compares well to the RAPID observations,
161 with a standard deviation value of 3.2 and 3.4 Sv, respectively, for the modeled and observed
162 monthly mean transports in 2004-2015 (Figure 4b). The magnitude of the modeled
163 interannual variability is slightly lower than observed, with a standard deviation of 1.0 Sv
164 versus 1.5 Sv for the annual mean transports (Figure 4b). On seasonal timescales, the
165 modeled AMOC variability is similar to observations in both magnitude (multi-year averaged
166 monthly mean transports give a standard deviation of 1.6 Sv) and phase, with low transports
167 in January-June and high transports in July-December (Figure 4c).

168 **3. Modeled and observed circulation in the South Atlantic**

169 In this section, we first compare the large-scale surface circulation to observations, and then
170 examine in more detail the transport structure along three sections in the South Atlantic:
171 34°S, 65°W in the Drake Passage, and a Prime Meridian-Good Hope section southwest of
172 Africa (Figure 1). Significant observations have been conducted at these locations and they
173 provide an important benchmark for evaluating the realism of the modeled transports,
174 which are used to document the transport structure of the South Atlantic.

175 **3.1 The surface circulation pattern**

176 Figure 5 compares the observed and modeled mean sea surface height (SSH), SSH variability,
177 and eddy kinetic energy (EKE) of the surface currents in the South Atlantic. The observed
178 mean SSH (Figure 5a) is from the climatology CNES-CLS13 (Rio et al., 2014) while the SSH
179 variability (Figure 5c) and surface EKE (Figure 5e) are derived from the AVISO data over
180 1993-2012, the same time period used for model results. In the western side of the domain,
181 part of the Antarctic Circumpolar Current (ACC) turns north after passing the Drake Passage
182 and becomes the Malvinas Current (also called the Falkland Current). The latter continues to
183 flow northward along the continental shelf of Argentina until it meets the southward flowing
184 Brazil Current south of the Rio de la Plata estuary near 36°S. The confluence of these two
185 western boundary currents with opposite directions and very different properties (warm
186 salty subtropical water versus cold fresh subantarctic water) leads to numerous high-energy
187 eddies and thus strong variability in this so-called Brazil-Malvinas confluence zone (Figures
188 5c-f). In the south, the ACC is mostly zonal and exhibits a contracted (stronger) front in two
189 areas: one near 40°W south of the Zapiola Drift and the other near 10°W over the MAR. We
190 note that the model results are in good agreement with the observations.

191 West of Africa, the numerical model exhibits a tongue of high SSH variability/EKE that
192 extends farther into the South Atlantic than is seen in observations. This is a common feature
193 for many eddying models (e.g., Maltrud and McClean, 2005; Dong et al. 2011) where the
194 Agulhas rings that shed from the Agulhas retroflection and translate northwestward into the
195 South Atlantic follow a regular pathway and are too energetic. Figure 6 further illustrates
196 this, showing the SSH variability for both the model and the observations along the Prime
197 Meridian over the 10-year period of 2003-2012. The figure shows that the modeled rings are
198 stronger and pass this longitude within a smaller latitudinal range than is seen in the
199 observations.

200 **3.2 The Meridional Circulation across 34°S**

201 To discuss the circulation/transport structure, we first examine the distribution of key water
202 masses in this region. Figure 7 displays a vertical section of the mean potential temperature
203 θ and salinity S at 34°S. The observations are based on the gridded monthly Argo profiles
204 (2004-2014) for the upper 2000 m and the World Ocean Atlas 2013 (WOA13, Locarnini et
205 al. 2013; Zweng et al. 2013) below 2000 m; the model results are 35-year means (1981-
206 2015). The water column at this latitude can be divided into four density layers of water
207 masses that are characterized most clearly in salinity (θ decreases monotonically): saline
208 near surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$), fresh Antarctic Intermediate Water (AAIW,
209 $35.65 < \sigma_2 < 36.58$), saline NADW ($36.58 < \sigma_2 < 37.12$), and fresh Antarctic Bottom Water
210 (AABW, $\sigma_2 > 37.12$). There is a general agreement in the temperature and salinity
211 distributions, although the salinity values differ somewhat between the model and
212 observations (e.g., fresher AAIW in the observations).

213 The time mean meridional velocity across 34°S and the corresponding volume transports for
214 the four water masses defined above are shown in Figure 8. The observations consist of
215 geostrophic transports derived from θ/S profiles (Argo-WOA13 data) and Ekman transports
216 from the wind stress; see Dong et al. (2014) for details. The model results are 35-year means
217 (1981-2015). The main circulation at this latitude consists of the South Atlantic subtropical
218 gyre (southward Brazil Current near the western boundary and northward interior flow)
219 and the AMOC (northward Bengula Current near the eastern boundary and southward
220 DWBC near the western boundary). Quantitatively, the modeled southward western
221 boundary current is about 50 Sv (15, 10, and 25 Sv for the surface water, AAIW, and NADW,
222 respectively), compared to 45 Sv observed (7, 8, and 30 Sv for the surface water, AAIW, and
223 NADW, respectively). In the observations, the subtropical gyre in the surface water and
224 AAIW layers extends from the western boundary to $\sim 10^\circ\text{E}$, while the northward flow of the
225 AMOC component occupies the rest of the section from $\sim 10^\circ\text{E}$ to the coast of Africa. The
226 modeled transport pattern is similar to the observations, except that the regular pathway of
227 the Agulhas rings leads to a north/south circulation in the Cape Basin. In the NADW layer,
228 both observations and model results show a strong southward DWBC west of 40°W and a
229 northward return flow east of 40°W . Note that the DWBC is quite wide at this latitude and
230 that the transport obtained by Meinen et al. (2017), 15 Sv west of 44.5°W , does not include
231 the full DWBC (30 Sv in Argo-WOA13 based observations and 25 Sv in model). The return
232 flow is distributed over a wide region between 40°W and 0°E in model whereas it is more
233 localized over the Walvis Ridge in the observations. In the Cape Basin, both observations and
234 model results show a recirculation of the NADW. This recirculation is likely driven by eddy
235 activity in the upper ocean and is stronger in the model (see Figure 5). The modeled AABW

236 transport is 1 Sv in the western basin, much less than the 4-7 Sv estimated in observations
237 (e.g., Hogg et al., 1982; Speer and Zenk, 1993). There is no northward AABW transport in the
238 Argo-WOA13 based results.

239 The meridional flows in Figure 8 show significant barotropic components, and the baroclinic
240 nature of the AMOC, i.e., northward flows in the upper limb and southward flows in the lower
241 limb, becomes apparent only when integrated across the basin (Figure 9). The zonally-
242 integrated mean transport stream functions with respect to depth z show a maximum
243 overturning depth of ~ 1300 m in both observations and model results (Figure 9a). The
244 modeled mean AMOC transport is 14 Sv. This value is in good agreement with the latest
245 estimate based on six years of moored observations at the western and eastern boundaries
246 (14.7 Sv, Meinen et al., 2018), but is significantly lower than the estimates based on XBT
247 transects (18 Sv, Dong et al., 2009; Garzoli et al., 2013) and Argo-WOA13 (20 Sv, Dong et al.,
248 2014). With respect to density (Figure 9b), the northward limb is above the density surface
249 (σ_2) 36.58 kg/m^3 and the southward limb below. The modeled mean AMOC transport is 15.3
250 Sv, compared to 18.7 Sv based on Argo-WOA13 (18.1 Sv based on XBT transects). The
251 modeled northward limb consists of 9.6 Sv of warm surface water transport and 5.7 Sv of
252 AAIW transport, compared to 12.7 and 6.0 Sv, respectively, in the observations. This implies
253 a lower meridional heat transport (MHT) of 0.35 ± 0.23 PW in the model, compared to
254 0.68 ± 0.24 PW in the Argo-WOA13 results. The historical estimates of the MHT near this
255 latitude are 0.22-0.62 PW (see Table 29.3 in Macdonald and Baringer, 2013).

256 The modeled AMOC transport variability is also lower than in the Argo-WOA13 based
257 observations, on both interannual and seasonal timescales (Figure 10): In 2004-2014, the

258 modeled annual mean transports have a standard deviation of 0.7 Sv, compared to 2.0 Sv in
259 observations (Figure 10a); the 11-year averaged monthly mean transports have a standard
260 deviation of 2.3 Sv in the model, compared to 3.3 Sv in observations (Figure 10b). Although
261 the magnitude is lower, the phase of the modeled seasonal variability is consistent with the
262 observations (Figure 10b).

263 **3.3 ACC Transport through the Drake Passage at 65°W**

264 The Drake Passage is the narrowest constriction of the ACC in the Southern Ocean and the
265 place where long-term sustained monitoring programs have been conducted; see Meredith
266 et al. (2011) for a review of historical observations. The canonical full-depth volume
267 transport is 133.8 ± 11.2 Sv, based on year-long mooring and cruise data obtained during the
268 International Southern Ocean Studies (ISOS, Whitworth, 1983; Whitworth and Peterson,
269 1985). Using a 20-year transport time series based on satellite altimetry data (1992–2012)
270 and in situ current meter data (2006–2009), this value was increased by Koenig et al. (2014)
271 to 141 Sv, with a standard deviation of 13 Sv and an error estimate of 2.7 Sv. More recently,
272 Chidichimo et al. (2014) and Donohue et al., (2016), estimated a significantly higher mean
273 ACC transport of 173.3 Sv (primarily due to a higher barotropic component) based on high-
274 resolution moored observations from 2007 to 2011.

275 In the model, the mean ACC transport is 155.5 Sv, in line with the observations. The modeled
276 mean zonal velocity through the Drake Passage at 65°W and the corresponding volume
277 transports for the four density layers defined earlier (surface water, AAIW, NADW, AABW)
278 are shown in Figure 11a. The ACC at this longitude exhibits four high velocity cores
279 (indicated by arrows in Figure 11a), corresponding to the ACC southern boundary (SBby,

280 south of 63°S), the southern ACC Front (SACCF, at 61-62°S), the Polar Front (PF, at 58-60°S),
281 and the Sub-Antarctic Front (SAF, at 55-55°S). These modeled fronts are at the same
282 locations as in the classical description of Orsi et al. (1995) based on hydrographic surveys
283 and as in the more recent results of Firing et al. (2011) using 4.5 years of shipboard ADCP
284 (acoustic Doppler current profiler) surveys; note the SAF is wide in Figure 11a because the
285 front turns northward at this longitude.

286 To examine the vertical structure, the modeled ACC transport at 65°W is shown in green in
287 Figure 12a. At that longitude, the modeled ACC transport decreases exponentially with
288 depth, as in the observations of Firing et al. (2011). In the top 1000 m, the modeled ACC
289 transport is 89 Sv, which is close to the 95 Sv in Firing et al. (2011) and the 90 Sv in Donohue
290 et al. (2016). One can further break down the transport into baroclinic and barotropic
291 components: at 65°W and above 3000 m, the modeled total, baroclinic, and barotropic
292 transports are 149.4, 112.7, and 36.7 Sv, respectively, compared to the observed 140, 112,
293 and 28 Sv in Koenig et al. (2014). Thus, the modeled ACC transport is higher than in Koenig
294 et al. (2014) because of a stronger barotropic component contribution. Note that the higher
295 total ACC transport obtained in Donohue et al. (2016) is also due to a higher barotropic
296 contribution (45.6 Sv). Overall, the barotropic component remains the main uncertainty in
297 determining the total ACC transport.

298 The modeled monthly mean and annual mean ACC transports have a standard deviation of
299 5.2 Sv and 2.1 Sv, respectively (Figure 13a). These numbers are relatively small compared to
300 the long-term mean value of ~155 Sv. The seasonal variability is also small (with a standard
301 deviation of 1.3 Sv) and exhibits a biannual pattern of high transports in April and October

302 and low transports in June and January (Figure 13b). These results agree with Koenig et al.
303 (2016).

304 **3.4 Zonal transports south of Africa**

305 The wide ocean gap between the Antarctica and southern tip of Africa makes it difficult to
306 fully measure the transport and its spatial structure. Observations have mostly focused on
307 measurements along the Prime Meridian (e.g., Whitworth and Nowlin 1987; Klatt et al. 2005)
308 to approximately 50°S and the Good Hope line from 0°E, 50°S to the Cape of Good Hope,
309 South Africa (e.g., Legeais et al. 2005; Gladyshev et al. 2008; Swart et al. 2008). We refer to
310 the combination of these two sections as the Prime Meridian-Good Hope (PM-GH) transect
311 (Figure 1), along which the velocity/transport structure is examined. The modeled net
312 transport through PM-GH (156.5 Sv) is essentially the same as the net transport through the
313 Drake Passage because of mass conservation, except for an additional 1 Sv from the Pacific-
314 to-Atlantic Bering Strait throughflow.

315 The circulation along the PM-GH section can be divided into three regimes (Figure 11b):

316 i) Weddell gyre south of 55.5°S. There are two eastward and two westward jets that form
317 the Weddell gyre. The two westward jets are found along the Antarctic Slope and the Maud
318 Rise (MR) near 64°S, whereas the two eastward jets are found near 58-59°S and along the
319 southern boundary (SBdy) of the ACC at 55.5°S right south of the Southwest Indian Ridge
320 (SIR). This jet pattern is consistent with the observations of Klatt et al., (2005, their Figures
321 4-5). The time mean transport of the modeled Weddell gyre is 47.1 Sv, compared to 55.7 Sv
322 in Klatt et al. (2005).

323 ii) ACC from 55.5 to 40°S. The modeled ACC exhibits high-velocity cores associated with the
324 SACCF (52°S), PF (50.4°S and 48°S), SAF (44.6°S), and the subtropical front (STF, 42°S)
325 respectively. These front positions are close to the observations based on repeat CTD/XBT
326 transects in this region (Swart et al. 2008, their Table 3). Note that the PF at this location is
327 split into two fronts, with the elevated eastward velocity between 47°S and 49°S
328 corresponding to its northern expression (Swart et al. 2008; Gladyshev et al. 2008). The
329 modeled STF is much weaker than any of the other ACC fronts as in the observations. The
330 modeled mean ACC transport across the PM-GH transect, defined as the transport from 55.5
331 to 40°S including the STF as in Orsi et al., (1995), is 177 Sv, compared to 147-162 Sv
332 estimated from CTD transects (Whitworth and Nowlin, 1987, Legeais et al., 2005, Gladyshev
333 et al., 2008). The modeled baroclinic transport is 106.4 Sv above 2500 m, compared to 84.7-
334 97.5 Sv derived from repeated hydrographic surveys and in combination with satellite
335 altimetry data (Legeais et al. 2005; Swart et al. 2008).

336 iii) Agulhas retroflexion and leakage north of 40°S. Farther north, the model results show a
337 pair of eastward and westward flows associated with the Agulhas retroflexion and Agulhas
338 Current. The 'net' transports north of 40°S is 12.0 and 9.4 Sv westward for the surface water
339 and AAIW. Thus, the Agulhas leakage provides more transport than the upper AMOC (~ 7 Sv
340 return eastward in the STF).

341 The modeled transport across the full PM-GH transect decreases with depth and is eastward
342 above 4000 m (blue line in Figure 12a). There is a weak westward flow below 4000 m. When
343 compared to the vertical structure of the transport in the Drake Passage (green line in Figure
344 12a), the eastward transport through PM-GH transect is weaker in the 0-600 m range and

345 stronger between 600 and 3500 m. The latter is due, in part, to the contribution of NADW
346 transport from the lower limb of the AMOC (red line in Figure 12a).

347 The modeled net transports into and out of the region bounded by a) 34°S, b), Drake Passage,
348 and c) PM-GH (see Figure1) is shown in Figure 12b. Between 1500 m and the bottom, there
349 is a net positive inflow. This implies a strong upwelling across 1500 m with a maximum
350 upward transport of 8.8 Sv. This is consistent with the picture put forward by Schmitz (1995)
351 and Talley (2013) that the Southern Ocean is a key upwelling region for NADW.

352 **4. Large-scale time mean circulation in the South Atlantic Ocean**

353 In the previous section, we showed that the model is able to represent the basic circulation
354 features of the South Atlantic and the Southern Ocean with reasonable accuracy. In this
355 section, we use the model results to further explore the time-mean circulation in the South
356 Atlantic. In particular, we focus on a) the diapycnal water mass transformations associated
357 with the upwelling as shown in the previous section, and b) the lateral circulation patterns
358 of the upper and lower limbs of the AMOC.

359 **4.1 Diapycnal water mass transformations**

360 Similar to Figure 12b, Figure 14a shows the net transports into and out of the region
361 bounded by the 34°S, 65°W, and PM-GH sections, but with respect to density layers. The
362 positive transports (black) denote water flowing into the region which has to exit into
363 another density layer flowing out of the region (white). A downward integration of these
364 layered transports gives the net transport between the sea surface and a given density
365 surface, and the difference between this net transport and the change of the volume above
366 the density surface over time gives the total diapycnal water mass transformation (black line

367 in Figure 14b) taking place in the region across that density surface; see Xu et al. (2018) for
368 a more detailed discussion on water mass transformation.

369 The results (black line in Figure 14b) show that there are transformations of intermediate
370 water (34.50-35.92 range) in this region: ~ 10 Sv from denser water across the density
371 surface of 35.92 and ~ 6 Sv from lighter water across 34.50. The diapycnal transformation
372 can be compared to the transformation (dashed black line in Figure 14b) that is driven
373 directly by the surface buoyancy fluxes, calculated from the surface density fluxes and
374 surface density using the thermodynamic method (e.g., Walin 1982; Speer and Tziperman,
375 1992; Brambilla et al., 2008; Langehaug et al., 2012; Xu et al., 2018). The result suggests that
376 the surface-forced transformation accounts for most of the area-integrated diapycnal
377 transformation in the South Atlantic region bounded by the three sections. Spatial
378 distribution of the transformation (not shown) further supports the notion that the
379 diapycnal transformation in this region is mostly due to surface buoyancy forcing: the ocean
380 loss buoyancy in the Agulhas Leakage region and the Brazil-Malvinas confluence zone, and
381 gain buoyancy in the ACC.

382 **4.2 AMOC horizontal circulation patterns**

383 **a) Upper limb (surface water and AAIW)**

384 The upper (northward) limb of the AMOC consists of two density layers: the surface water
385 ($\sigma_2 < 35.65$) and the AAIW ($35.65 < \sigma_2 < 36.58 \text{ kg m}^{-3}$). The modeled 35-year (1981-2015) mean
386 horizontal circulation for these two layers is displayed in Figures 15 and 16, respectively.
387 There are three distinct pathways for the surface water (Figure 15): first, the AMOC
388 component which flows directly northwestward (the so-called Agulhas Leakage) into the

389 South Atlantic (red stream lines); second, the subtropical ‘supergyre’ that flows around the
390 subtropical gyre and connects the South Atlantic and Indian Ocean (De Ruijter, 1982; Gordon
391 1992) (blue lines); and third, the subtropical gyre (orange lines). There is virtually no surface
392 water in the ACC from the Pacific Ocean (pink lines). In addition, this surface water is
393 prevented from moving north in the South Atlantic because of the ‘supergyre’ and cannot
394 contribute directly to the AMOC.

395 The circulation pattern of the modeled AAIW (Figure 16) is similar to the surface water
396 (Figure 15), but it shows a meridionally more confined subtropical gyre (orange lines) and a
397 larger contribution to the ACC from the Pacific Ocean (pink lines). As for the surface water,
398 there is a ‘supergyre’ connecting the subtropical gyres of the South Atlantic and Indian
399 Oceans, which prevents a direct contribution of water mass from the ACC into the upper limb
400 of the AMOC. The modeled circulation patterns of Figures 15 and 16 are similar to the
401 schematic of Stramma and England (1999, their Figures 3-4) except for the recirculation in
402 the Cape Basin, which is due to the unrealistic pathways of the modeled Agulhas eddies (see
403 Figures 5 and 6 and discussion in section 3.1).

404 Although the AAIW in the ACC does not directly contribute to the upper limb of the AMOC,
405 the cold Pacific water entering the South Atlantic via the Drake Passage does play a role in
406 setting the subtropical gyre water mass characteristics. To illustrate this, we project the
407 northward volume transports across the 45°S, the 34°S, and the GH sections (Figure 16) in
408 terms of potential temperature and salinity (θ -S). The results show a) the Pacific AAIW
409 (35.65-36.58) that flows northward across 45°S (Figure 17a) is much colder and fresher than
410 the Indian AAIW that flows westward across the GH section (Figure 17c); and b) the AAIW

411 that flows northward across 34°S (Figure 17b) is a combination of these two water masses.
412 About 9.2 Sv of the northward AAIW transport across 34°S exhibit properties similar to the
413 warmer saltier Indian water at the GH section, whereas about 8.3 Sv exhibit properties
414 similar to the colder fresher Pacific water at 45°S. Thus, although the cold Pacific water does
415 not directly feed into the upper limb of the AMOC, it contributes to the water properties at
416 34°S and impacts the meridional heat/freshwater transports.

417 **b) Lower limb (NADW)**

418 Figure 18 shows the modeled mean circulation for the NADW density layer
419 ($36.58 < \sigma_2 < 37.12$). The modeled NADW flows southward as a DWBC along the continental
420 slope of the Brazil and the Argentine Basins, all the way to about 40°S where it encounters
421 the northward-flowing deep Falkland Current. The NADW continues to flow southward (now
422 offshore of the deep Falkland Current) to about 45°S where it meanders and flows eastward
423 south of the Zapiola Drift (Rise). This modeled NADW pathway is similar to the one described
424 in the schematic of Stramma and England (1999, their Figure 5) and is consistent with
425 picture derived from salinity, oxygen, and other tracers (CFC, for example) (e.g., Koltermann
426 et al., 2011; Garzoli et al., 2015). There is a strong counter-clockwise flow around the Zapiola
427 Drift (Figure 18). Within the NADW layer across 45°S, the modeled transport is about 45 Sv
428 for the southward flow in the western side and 25 Sv for northward flow in the eastern side.
429 For the full water column, the modeled long-term mean transport for this Zapiola anticyclone
430 is on the order of 90 Sv in the western side and 50 Sv in the eastern side. Saunders and King
431 (1995) proposed a flux of 80 Sv based on one CTD/ADCP survey along 45°S, but they

432 cautioned that the feature is “unexpected and not entirely free from uncertainty.” The model
433 results suggest that their estimate actually may be quite reasonable.

434 Offshore of the DWBC, there are complex recirculation patterns in the Brazil Basin, especially
435 around the Vitoria-Trindade Seamount Chain near 20°S (Hogg and Owens, 1999). The
436 recirculation carries NADW from the DWBC toward the interior of the basin and explains the
437 high salinity signature found along the WOCE lines A09 and A095 that extend all the way to
438 the MAR in both the observations and the model (Figure 19). The model results exhibit a
439 zonal flow of about 2 Sv across the MAR near 22°S (see Figure 20 over the bathymetry). This
440 zonal flow has been the subject of a number of studies, with an estimated transport ranging
441 from 2.2 to 5 Sv (i.e., Warren and Speer, 1991; Speer et al., 1995; Hogg and Thurnherr, 2005;
442 Garzoli et al., 2015) although 10.7 Sv in Arhan et al. (2003). Note in Figure 19 that the salinity
443 is similar between A15 and A16, but quite different between A15 and A14. The salinity
444 contrast does not support a high transport across the MAR.

445 South of 22°S, there are weaker, more widespread westward zonal flows across the MAR
446 (Figures 18 and 20). As a consequence, the interior salinity in the west basin near 30°S (along
447 A10) is much lower compared to 20-25°S (A09 and A095). Also, the modeled net transport
448 across the MAR is close to zero and virtually all the NADW flows across the MAR south of
449 40°S. As the NADW from north and the ACC water from the Drake Passage flow eastward,
450 the streamlines turn northward when approaching the MAR and southward after crossing
451 the MAR. This meridional shift can be explained by the conservation of potential vorticity,
452 i.e., a decrease in thickness h when approaching the MAR leads to a northward shift to reduce
453 the planetary rotation f so that f/h is constant and vice versa. Because the MAR is slanted in

454 a northwest-to-southeast direction in this area, the northward and southward shifts at
455 different longitude lead to a contraction of the front near 10°W , clearly visible in the SSH for
456 both model and observations (Figure 5).

457 **5. Latitudinal coherence of the AMOC variability**

458 In this section, we examine whether the modeled AMOC variability is latitudinally coherent
459 throughout the Atlantic from 35°S to 70°N . Using hydrography and satellite data with a box
460 model, Kelly et al. (2014) suggest that the meridional heat transport anomalies (closely
461 correlated with the strength of the AMOC) are highly coherent from 35°S to 40°N on
462 interannual timescales. Xu et al. (2014) found a similar result, but their model was regional
463 (did not extend past 25°S) and was only integrated for eight years after spin up. The model
464 used in this study is global and was integrated for close to six decades. To examine the
465 latitudinal coherence on different timescales, we focus on the last 35 years (1981-2015) of
466 the simulation to ensure that the solution is fully spun up. First, we compute monthly mean
467 AMOC transports at all latitudes from 35°S to 70°N (in 0.5° increments) and then decompose
468 the transport time series into different timescales using the ensemble empirical mode
469 decomposition (EEMD, Wu and Huang, 2009). The EEMD extracts the amplitude-frequency
470 modulated oscillatory components (termed as “intrinsic mode functions”, or IMFs)
471 successively from the highest to the lowest frequencies, without using a priori determined
472 basis function. The advantage of this method is that both the frequency and amplitude of
473 each IMF are determined adaptively from the local characteristic of the time series and vary
474 as a function of time. Finally, the same IMF at each latitude is pieced together to construct
475 the basin-scale AMOC variability.

476 The number of IMFs is the integer of the binary logarithm of the time series length N ($\log_2 N$),
477 or 7 for $N=420$ months. The first two IMFs represent the high-frequency intraseasonal
478 variability and are not shown. The third IMF (Figure 21) exhibits frequency close to seasonal
479 timescale. The results show that the seasonal AMOC variability is high, especially in the
480 equatorial region due to the seasonal migration of the Inter-tropical convergence zone
481 (ITCZ), and is coherent across different latitudes. To the south and north of about 20°N , there
482 is a shift in the phase of the variability. These results are very similar to the regional Atlantic
483 model of Xu et al. (2014). The phase shift seen in the model can also be seen in observations:
484 for example, high seasonal AMOC transports occur during April-August at 34°S (Figure 10b
485 based on Argo-WOA13 data), compared to July-November at 26.5°N (Figure 4c based on
486 RAPID data).

487 The fourth and fifth IMFs are presented in Figure 22 and they represent interannual
488 variability. The interannual variability is weaker than seasonal variability; it also exhibits a
489 good coherence throughout the Atlantic domain. Near 40°N , there is a phase shift with
490 variability in the subpolar North Atlantic leading the variability in the subtropics. This shift
491 is clearer in the fifth IMF (Figure 23b). South of 35°N , the interannual variability is very
492 similar in phase and is consistent with the finding of Kelly et al. (2014) based on
493 observations, i.e., that the interannual variability of the meridional heat transport (a good
494 index for AMOC) is highly coherent between 35°S and 40°N on these timescales.

495 The sixth and seventh IMFs (Figure 23) represent the AMOC variability on decadal (about
496 10-30 years) timescales. The decadal variability is even weaker than interannual timescales
497 and exhibits a good coherence. Similar to the interannual variability shown in Figure 22b,

498 the decadal variability in Figure 23a (10-year) shows a phase shift near 40°N with variability
499 of the subpolar North Atlantic leading variability farther south. The variability in Figure 23b
500 generally shows a high AMOC transport period from 1985-2000 and a lower AMOC transport
501 from 2000-2015, a close to 30-year variation that is coherent from 35°S to 50°N.

502 **6. Summary and Discussion**

503 Through the South Atlantic Ocean, the AMOC is connected to the Indian and Pacific Oceans
504 and is entangled into the global overturning circulation system. This important region is also
505 particularly complex, featuring strong boundary currents (jets) and high eddy variability in
506 both the western and eastern boundaries as well to the south in the Atlantic sector of the
507 Southern Ocean. Observations of the full-depth circulation structure are focused on limited
508 places, thus the 3-D circulation structure in the South Atlantic and the large-scale pattern of
509 the AMOC variability are not well-determined. In this study, we used numerical results from
510 a long-term 1/12° global simulation, along with observations, to address the fundamental
511 questions. The model results are shown to represent the transports and the vertical
512 structure of the key circulation patterns in this region, especially, the AMOC across 34°S in
513 the South Atlantic, the ACC at 65°W in the Drake Passage, as well as the zonal flows along the
514 PM-GH transect in the open ocean southwest of Africa. The key results are

515 1) Strong upwelling takes place in the South Atlantic Ocean. In the area south of 34°S
516 between the Drake Passage (65°W) and southwest of Africa (PM-GH), there is a maximum
517 upward transport of 8.8 Sv across 1500 m. The water in this area also undergoes a strong
518 diapycnal transformation that forms about 16 Sv intermediate water in 34.50-35.92 density
519 range: 10 Sv from denser AAIW and NADW and 6 Sv from lighter near surface water.

520 2) In the upper limb of the AMOC, the northward flow originates from the warm Agulhas
521 leakage. The existence of a super gyre (connecting the subtropical gyre of Atlantic and Indian
522 Oceans) prevents the colder fresher Pacific water from directly contributing to the AMOC.
523 The cold water, however, modifies the water properties of both the AMOC and the
524 subtropical gyre across 34°S, thus contributing to the meridional transports of heat and
525 freshwater (salinity anomaly).

526 3) In the lower limb, the NADW flows southward in the DWBC along the continental slope
527 and in complex recirculation in the Brazil Basin, especially around the Vitória-Trindade
528 Seamount Chain near 20°S. The recirculation carries the NADW and its high-salinity
529 signature into the offshore interior. A zonal flow of NADW crossing the MAR is found near
530 22°S. This zonal flow is likely minor given the fact that the salinity/oxygen east of MAR is
531 significantly lower as compared to the west. Virtually all of the NADW from the north flows
532 in the DWBC all the way to 40-45°S before turning eastward to flow across the MAR near
533 42°S. The location of this crossing can be seen as a contracted ACC front near 10°W.

534 4) Overall, the modeled AMOC variability from seasonal to decadal timescales (Figure 21-23)
535 shows a good meridional coherence throughout the Atlantic Ocean, especially from 35°S to
536 about 35°N, consistent with the finding by Kelly et al. (2014). The coherence in this latitude
537 range, where the diapycnal water mass transformation between the upper and lower AMOC
538 limbs is expected to be small, indicates that the AMOC variability is modulated by deep water
539 formation in the subpolar North Atlantic and/or upwelling regions. The decadal variability
540 shows some phase shift near 40°N, with variability of the subpolar North Atlantic leading
541 that to the south. A similar phase shift near this latitude is found in the coupled climate

542 simulation of Zhang (2010), but not robust across in the coupled CMIP5 and in the
543 atmospherically-forced global ocean-sea ice simulations at similar resolution, i.e., the CORE-
544 II (Xu et al. 2018b).

545 Some model features need to be improved. One key challenge for numerical models in this
546 region is to represent the northwestward translation of the Agulhas rings into the South
547 Atlantic. Compared to observations, the modeled rings are dissipated slower and follow a
548 more regular pathway. This leads to a high EKE tongue extending much farther to the
549 northwest and has significant impact on the regional circulation pattern in the eastern South
550 Atlantic. Possible improvements may result from a) including the ocean currents in the wind
551 stress calculation, which adds a significant damping effect to eddies (Renault et al., 2017);
552 and b) using a higher order advection scheme in momentum equation, which leads to more
553 irregularity in eddy size and pathway (Backeberg et al., 2009).

554 Furthermore, the results presented in this study are based on a single simulation so the
555 robustness needs to be examined with coordinated model comparison efforts, with the same
556 model resolution and atmospheric forcing. Such multi-model inter-comparison efforts have
557 been shown to be useful in the low-resolution, atmospherically-forced global ocean-sea ice
558 simulations under the CORE-II and the fully-coupled climate simulations under CMIP5 (e.g.,
559 Xu et al., 2018b). A similar effort with several ocean general circulation models (OGCMs) at
560 $1/12^\circ$ resolution, all using the latest atmospheric forcing JRA-55 (Tsujino et al., 2018), is
561 ongoing and the robustness of the modeled circulation in this important region will be
562 further investigated.

563

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574 measurements are available at http://www.argo.ucsd.edu/Gridded_fields.html and the
575 World Ocean Atlas 2013 is available at <http://www.nodc.noaa.gov/OC5/woa13>; the velocity
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577 global model outputs are stored in the U.S. Navy DSRC archive server and the model results
578 presented in this study are available in HYCOM server
579 (<ftp://ftp.hycom.org/pub/xbxu/GLBb0.08/SATL>).

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817

818 **Figure Captions:**

819 Figure 1: Model bathymetry (in km) along with key topographic features in the South Atlantic Ocean.
 820 Red lines denote three sections where significant observations have been obtained and the
 821 observations are used to evaluate the model results: 34°S in the South Atlantic, 65°W in the Drake
 822 Passage, and the Prime Meridian-Good Hope (PM-GH) transect southwest of Africa.

823 Figure 2: Time evolution of global domain-averaged a) potential temperature in °C and b) total kinetic
 824 energy in cm^2s^{-2} from the global ocean-sea ice simulation. Thin and thick lines denote monthly and
 825 annual means, respectively.

826 Figure 3: Time evolution of a) monthly mean sea ice extent and b) annual mean sea ice extent
 827 anomaly relative to the 1980-2015 average. The sea ice extent is defined as the area (in 10^6 km^2) with
 828 sea ice concentration of 15% or higher. The red/blue lines are model results in the
 829 northern/southern hemisphere; black lines are observations from NSIDC (National Snow and Ice
 830 Data Center).

831 Figure 4: a) Time evolution of the modeled monthly mean transport of the Atlantic meridional
 832 overturning circulation (AMOC) at 26.5°N; blue dash line denotes a decreasing trend of 0.43 Sv per
 833 decade; b) Variability of the AMOC transports during 2004-2015, with thin/thick lines denoting
 834 monthly/annual means; c) Seasonal variability of the AMOC transports at 26.5°N averaged over
 835 2004-2015. Red/black in panels b-c) are model results and observations from the RAPID array.

836 Figure 5: Observed and modeled distributions of a-b) time mean sea surface height (SSH, in cm), c-d)
 837 SSH standard deviation (in cm), and e-f) eddy kinetic energy (EKE, in $\text{cm}^2 \text{ s}^{-2}$) of the surface current
 838 in the southern Atlantic. In observation, the mean SSH is based on long-term climatology CNES-CLS13
 839 (Rio et al., 2014); the SSH standard deviation and EKE are based on AVISO data in 1993-2012. All
 840 model results are in 1993-2012.

841 Figure 6: Observed (upper panel) and modeled (lower panel) sea-surface height (SSH) anomaly along
 842 the Prime Meridian in latitude range 20-40_S from 2003 to 2012. The black '+' denotes the latitude
 843 of the Agulhas Rings that pass this longitude.

844 Figure 7: Observed and modeled potential temperature and salinity across 34°S. Observations are
 845 based on a combination of Argo profiles (2004-2014) for the top 2000m and World Ocean Atlas 2013
 846 (WOA13) below 2000 m; model results based on the global $1/12^\circ$ HYCOM simulation in 1981-2015;
 847 The three black lines denote isopycnic interfaces that divide the water column into four layers: near
 848 surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$), Antarctic Intermediate Water (AAIW, $35.65 < \sigma_2 < 36.58$), North
 849 Atlantic Deep Water (NADW, $36.58 < \sigma_2 < 37.12$), and Antarctic Bottom Water (AABW, $\sigma_2 > 37.12$).

850 Figure 8: Observed and modeled time mean meridional velocity across 34°S and the corresponding
 851 volume transport for the four density layers: surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$), Antarctic
 852 Intermediate Water (AAIW, $35.65 < \sigma_2 < 36.58$), North Atlantic Deep Water (NADW, $36.58 < \sigma_2 <$
 853 37.12), and Antarctic Bottom Water (AABW, $\sigma_2 > 37.12$). Observations based on a combination of
 854 Argo-WOA13 profiles; model results based on global $1/12^\circ$ HYCOM simulation in 1981-2015.

855 Figure 9: a) Long-term mean meridional overturning streamfunction (in Sv) at 34°S with respect to
 856 a) depth and b) potential density in σ_2 . Observations based on monthly mean Argo profile for the
 857 upper 2000m and WOA13 below 2000 m; model results based a global 1/12° HYCOM simulation
 858 (1981-2015).

859 Figure 10: a) Time series of the observed (black) and modeled (red) AMOC transport at 34°S, with
 860 thin/ thick lines denoting monthly/annual means, respectively; b) seasonal variability of the AMOC
 861 transports at 34°S, averaged over 2004-2014 in observations (black) and 1981-2015 in model
 862 results(red).

863 Figure 11: Modeled long-term mean zonal velocity and the corresponding four-layer volume
 864 transport in four density layers across a) 65°W in the Drake Passage and b) the Prime Meridian-Good
 865 Hope (PM-GH) transect southwest of Africa. The ∇ denote the locations of Antarctic circumpolar
 866 current (ACC) fronts, from south to north, the Southern Boundary, South ACC Front, Polar Front,
 867 Subantarctic Front, as well as the subtropical front (only in panel b). The shaded area in panel b)
 868 between 40 and 55.5°S marks the ACC regime across the PM-GH transect. Transport are accumulative
 869 northward. The four layers are near surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$), Antarctic Intermediate Water
 870 (AAIW, $35.65 < \sigma_2 < 36.58$), North Atlantic Deep Water (NADW, $36.58 < \sigma_2 < 37.12$), and Antarctic
 871 Bottom Water (AABW, $\sigma_2 > 37.12$).

872 Figure 12: a) Modeled mean horizontal transports (in Sv) for very 100 m in the vertical across the
 873 34°S, the 65°W, and the PM-GH transects. b) The net transports into the region enclosed by the three
 874 transects, with positive (negative) values indicating net transport into (out of) the region.

875 Figure 13: a) Variability of the modeled ACC transport through the Drake Passage at the 65°W, in
 876 monthly (thin) and annual (thick) means; b) Seasonal variability of the modeled ACC transports
 877 averaged over 1981-2015.

878 Figure 14: a) Net transports into the South Atlantic region closed by the 34°S, the 65°W, and the PM-
 879 GH transects, with respect to density layers (positive/negative values for net transport into/out of
 880 the region); b) Black line denotes the total diapycnal transformation; and dashed red line denotes the
 881 surface forced diapycnal transformation calculated from surface buoyancy fluxes.

882 Figure 15: Modeled long-term mean horizontal transport streamfunction (in Sv) for the layer of near
 883 surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$). Each streamline contour is 2 Sv. Red, blue, and orange streamlines
 884 denote AMOC contribution, super gyre that ow around the South Atlantic and back to Indian Ocean,
 885 and subtropical gyre of the South Atlantic.

886 Figure 16: Modeled long-term mean horizontal transport streamfunction (Sv) for the layer of AAIW
 887 ($35.65 < \sigma_2 < 36.58 \text{ kg m}^{-3}$). Thick pink stream lines (increment of 8 Sv) is the ACC. The Red, blue,
 888 and orange streamlines denote AMOC contribution, super gyre, and subtropical gyre of the South
 889 Atlantic (similar to Figure 15). The dashed blue lines denote 34°S, 45°S, and the GoodHope sections,
 890 across which the water properties of the northward and northwestward transports are examined
 891 Figure 17.

892 Figure 17: Modeled northward transport (in Sv) across 45°S and 34°S, and northwestward transport
893 across the GH section, projected on potential temperature-salinity (θ - S) plane with $\Delta\theta \times \Delta S$ of
894 $0.2^\circ\text{C} \times 0.04$. The isopycnal (σ_2) surfaces of 35.65 and 36.58 kg m^{-3} denote the upper and lower AAIW
895 interfaces.

896 Figure 18: Modeled long-term mean horizontal transport streamfunction for the layer of NADW
897 ($36.58 < \sigma_2 < 37.12 \text{ kg m}^{-3}$). Pink streamlines (10 Sv increment) indicate the eastward transport of
898 the ACC, blue to yellow streamlines (2 Sv increment) represent the southward spreading of the
899 NADW from north.

900 Figure 19: a) observed and model salinity distribution at 2500m in the South Atlantic. Observation
901 based on CTD data from GoShip program <http://www.go-ship.org>. Detailed vertical sections can be
902 seen in the WOCE Atlas (Kiltermann et al., 2011). The results show an eastward extension of high
903 salinity (NADW signature) between 20 and 25°S west of the mid-Atlantic Ridge (MAR), and
904 significantly lower salinity east of MAR.

905 Figure 20: Zoomed view of the circulation for the density layer of NADW ($36.58 < \sigma_2 < 37.12 \text{ kg m}^{-3}$)
906 across the Mid-Atlantic Ridge in the South Atlantic Ocean. The blue contours denote NADW from
907 north and magenta streamlines denote ACC from Drake Passage.

908 Figure 21: Modeled seasonal variability of the AMOC transports at different latitude, based on the
909 third Intrinsic mode function (IMF) using the ensemble empirical mode decomposition (EEMD,
910 Huang and Wu (2008) and Wu and Huang (2009).

911 Figure 22: Modeled interannual variability of the AMOC transports at different latitude, based on the
912 fourth and fifth Intrinsic mode function (IMF) using the ensemble empirical mode decomposition
913 (EEMD, Huang and Wu (2008) and Wu and Huang (2009).

914 Figure 23: Modeled decadal variability of the AMOC transports at different latitude, based on the sixth
915 and seventh Intrinsic mode function (IMF) using the ensemble empirical mode decomposition
916 (EEMD, Huang and Wu (2008) and Wu and Huang (2009).

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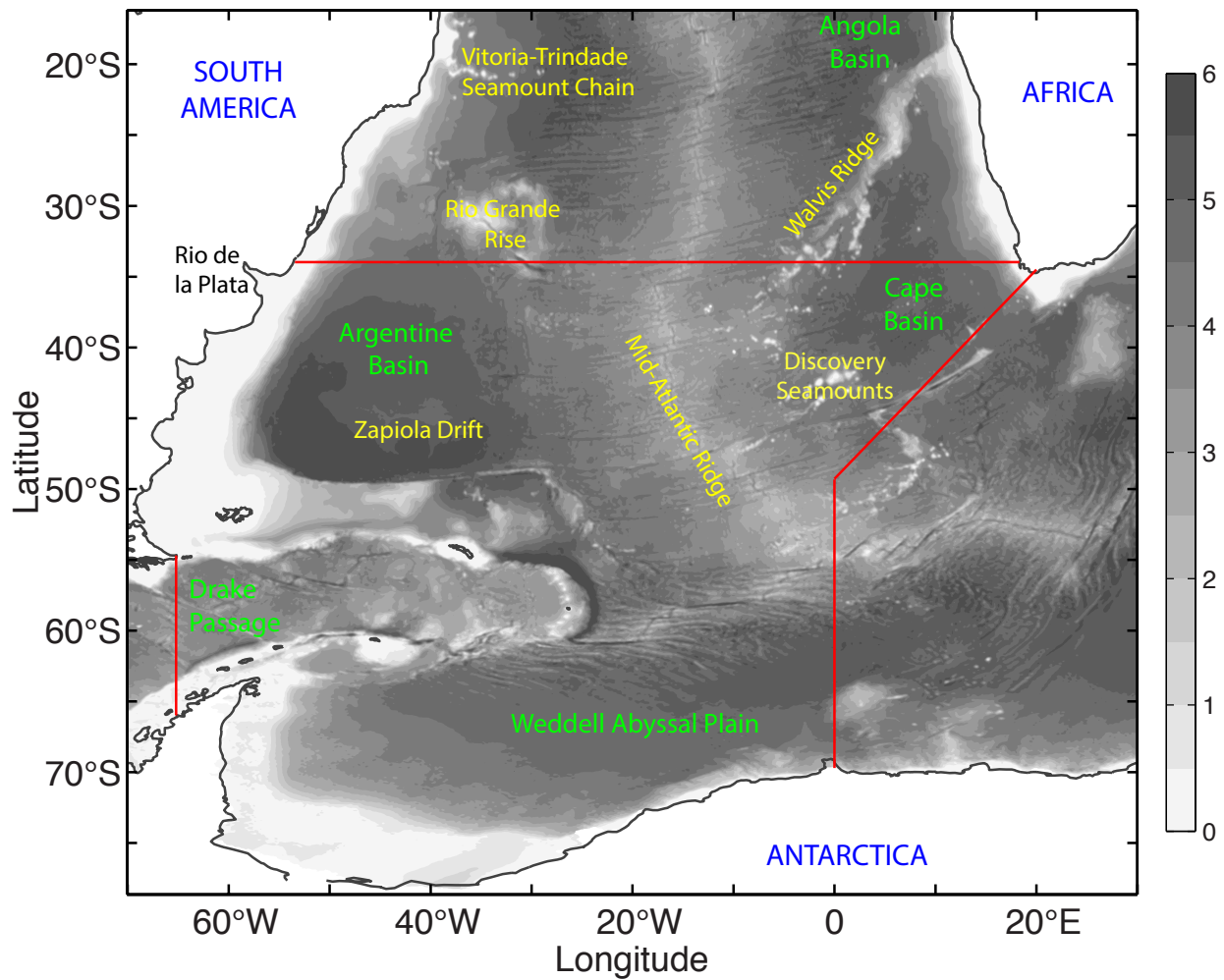


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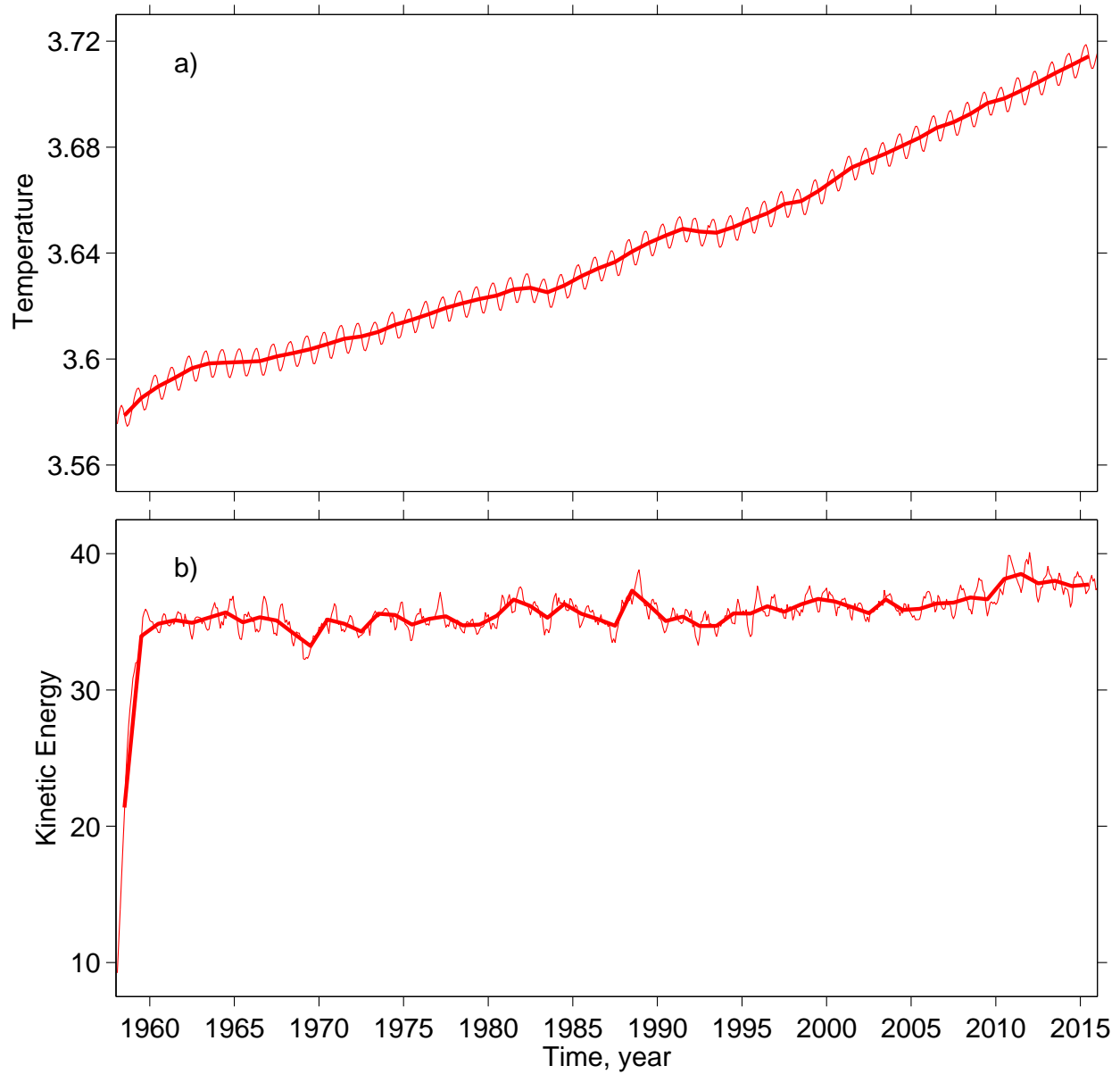


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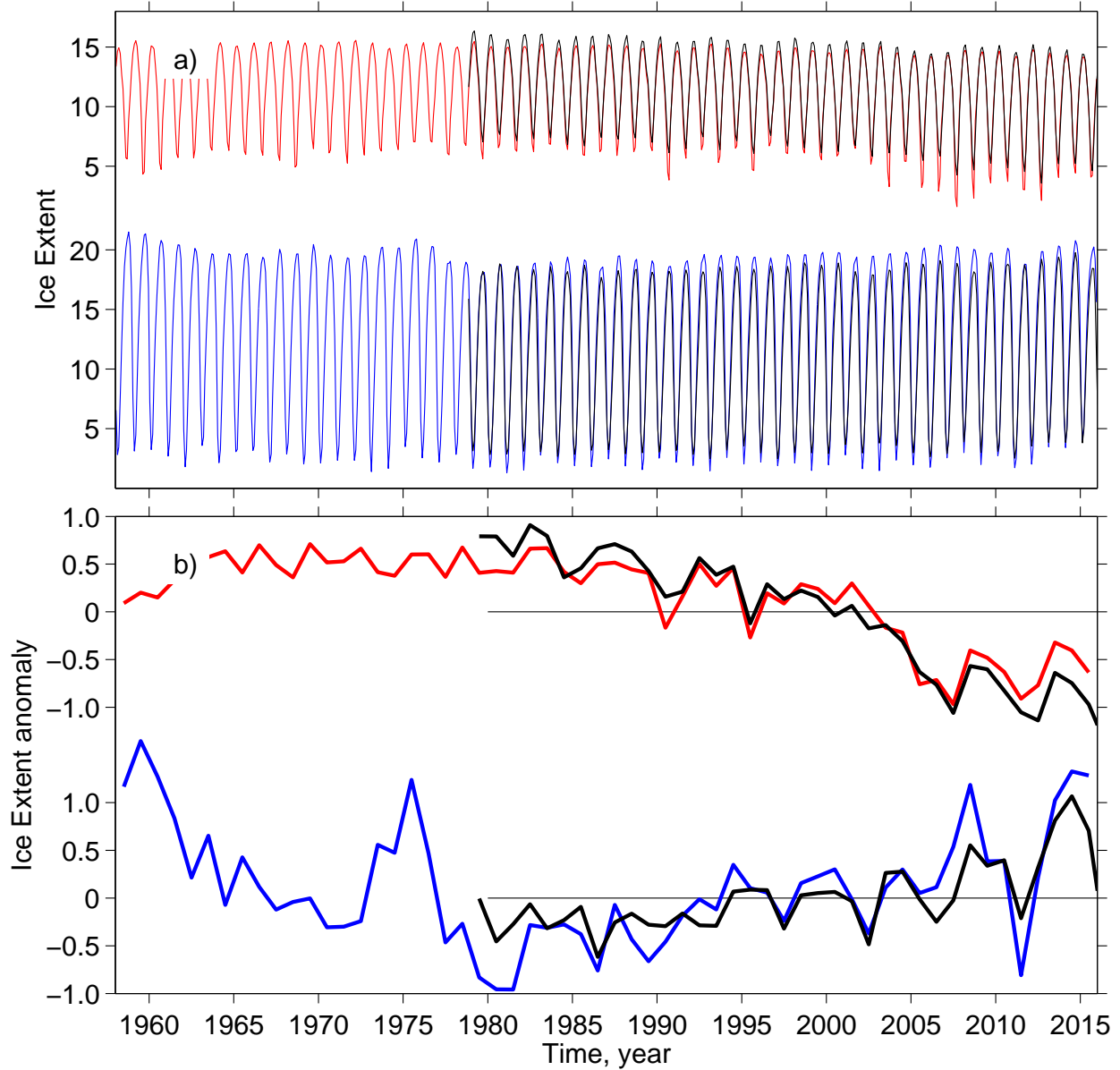


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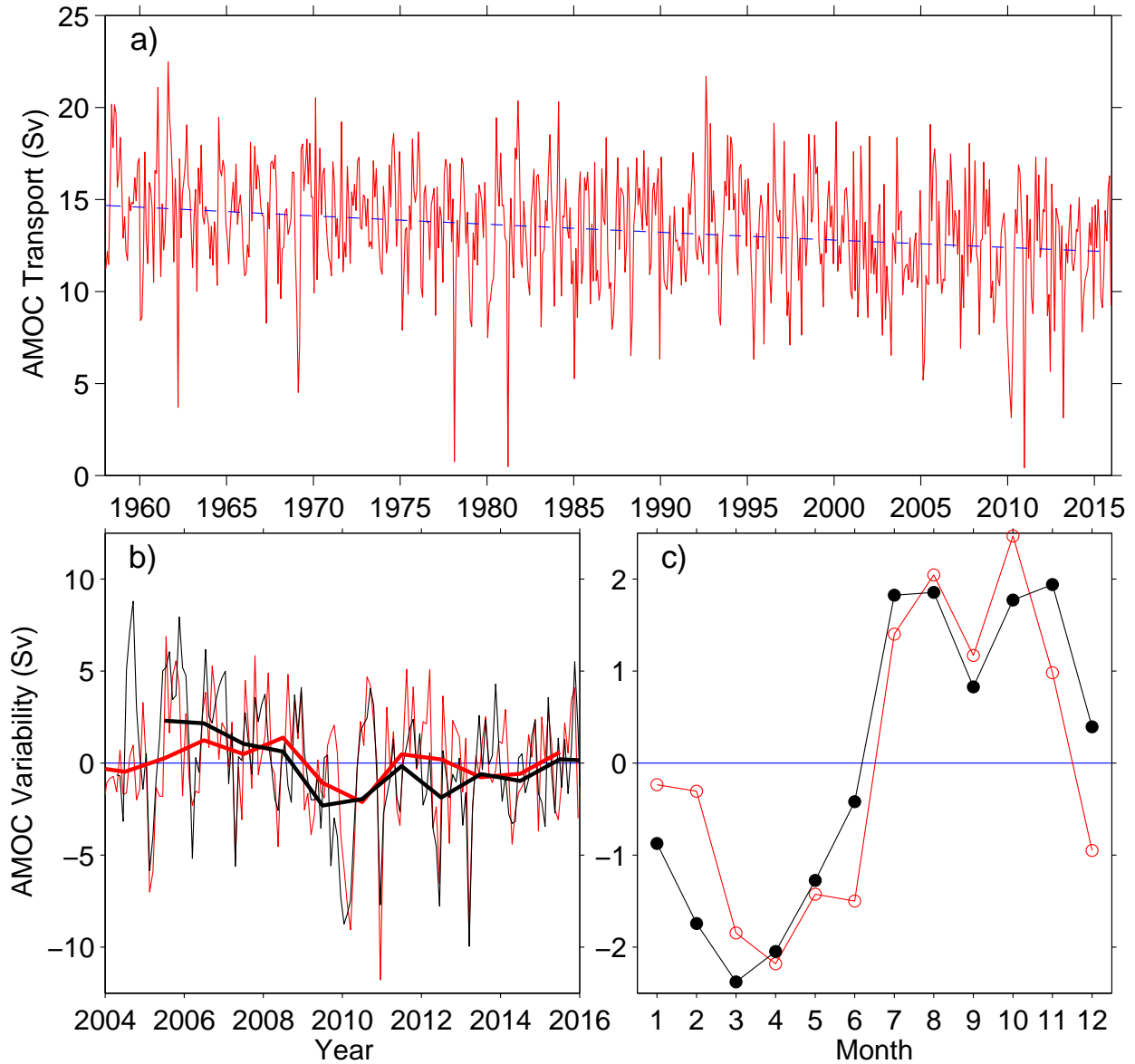


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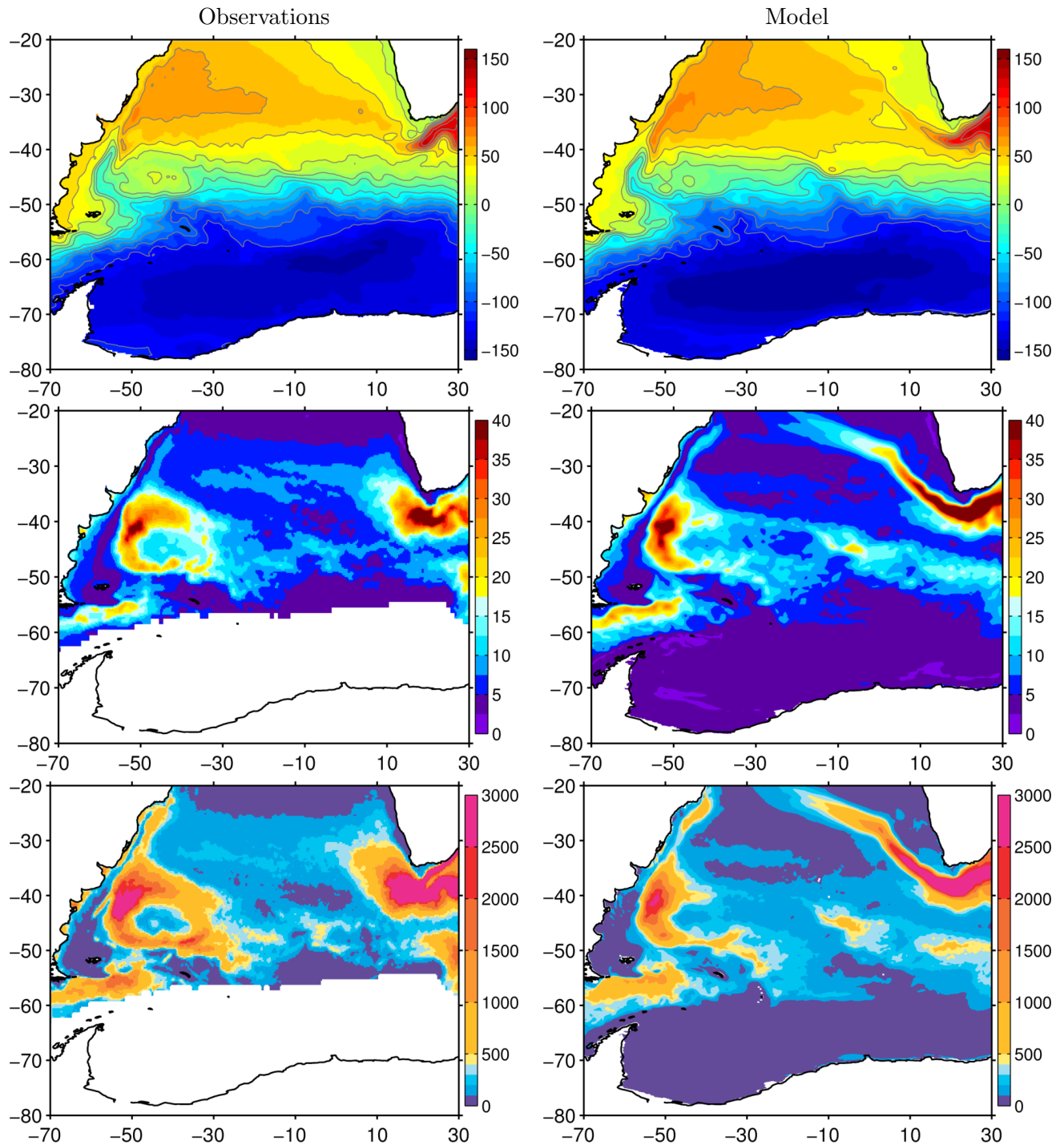


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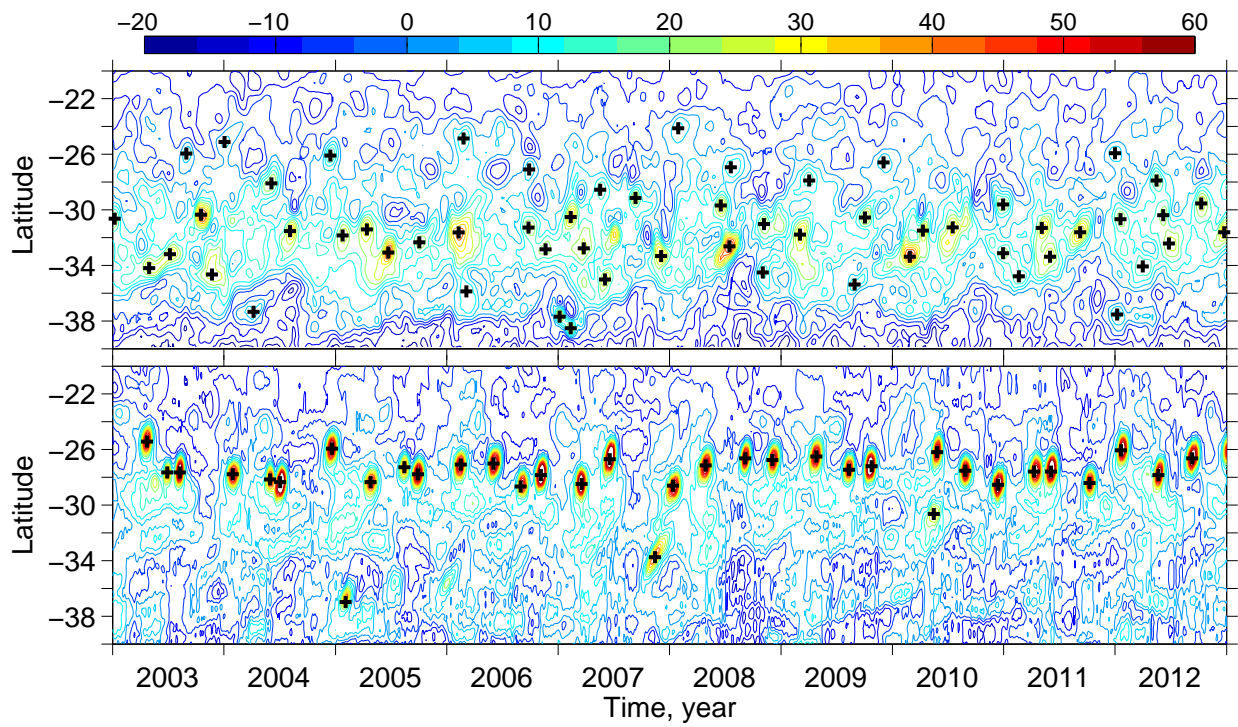


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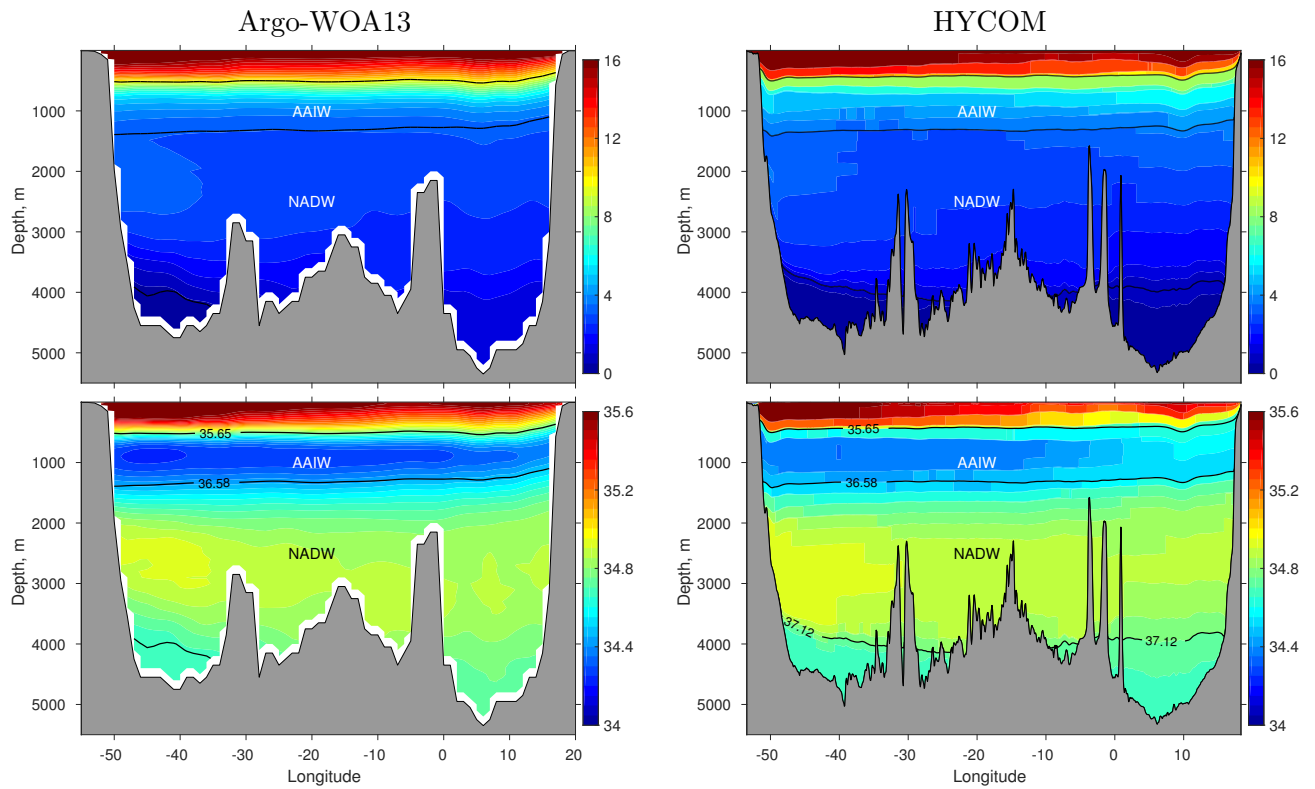


Figure 7: Observed and modeled potential temperature and salinity cross-sections across 34°S . Observations are based on a combination of Argo profiles (2004-2014) for the top 2000 m and World Ocean Atlas 2013 (WOA13) below 2000 m; model results based on the global $1/12^{\circ}$ HYCOM simulation in 1981-2015; The three black lines denote isopycnic interfaces that divide the water column into four layers: near surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$), Antarctic Intermediate Water (AAIW, $35.65 < \sigma_2 < 36.58$), North Atlantic Deep Water (NADW, $36.58 < \sigma_2 < 37.12$), and Antarctic Bottom Water (AABW, $\sigma_2 > 37.12$).

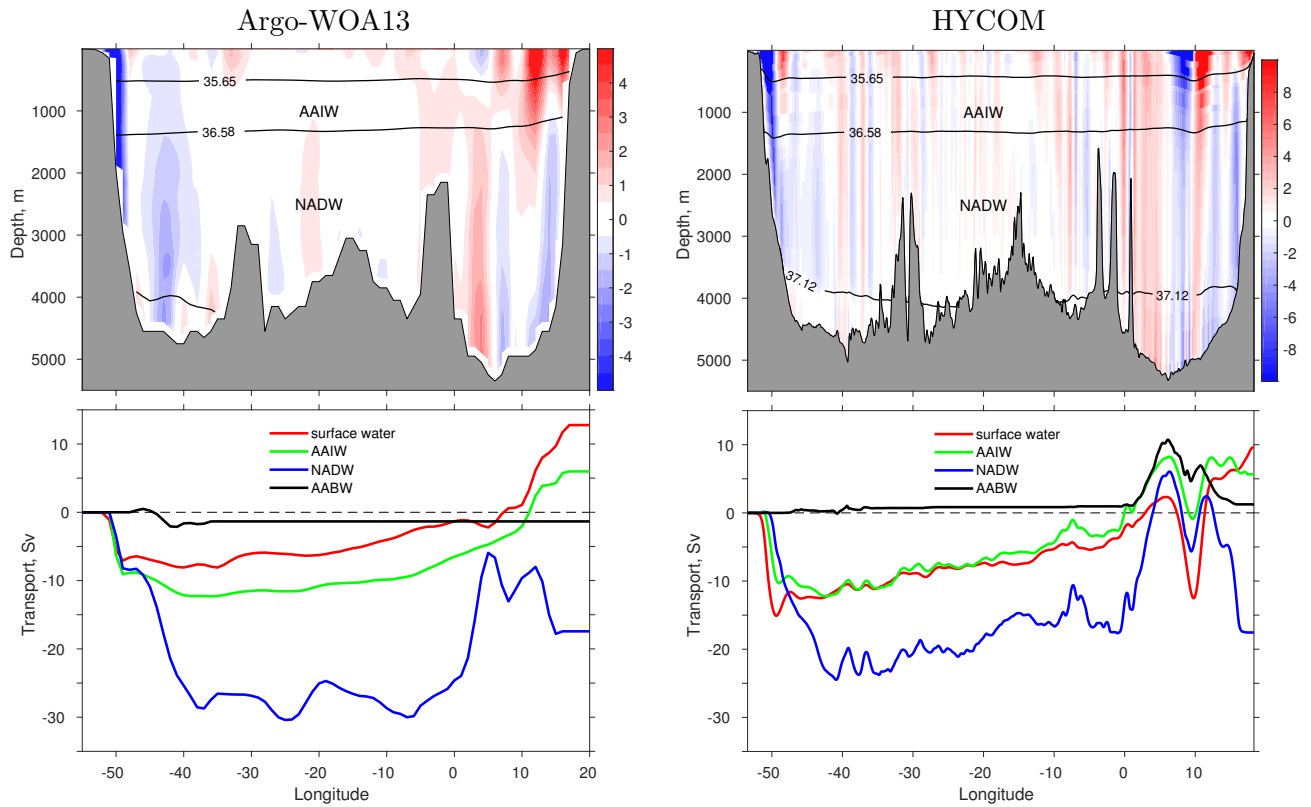


Figure 8: Observed and modeled time mean meridional velocity across 34°S and the corresponding volume transport for the four density layers: near surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$), Antarctic Intermediate Water (AAIW, $35.65 < \sigma_2 < 36.58$), North Atlantic Deep Water (NADW, $36.58 < \sigma_2 < 37.12$), and Antarctic Bottom Water (AABW, $\sigma_2 > 37.12$). Observations based on a combination of Argo-WOA13 profiles; model results based on global $1/12^{\circ}$ HYCOM simulation in 1981-2015.

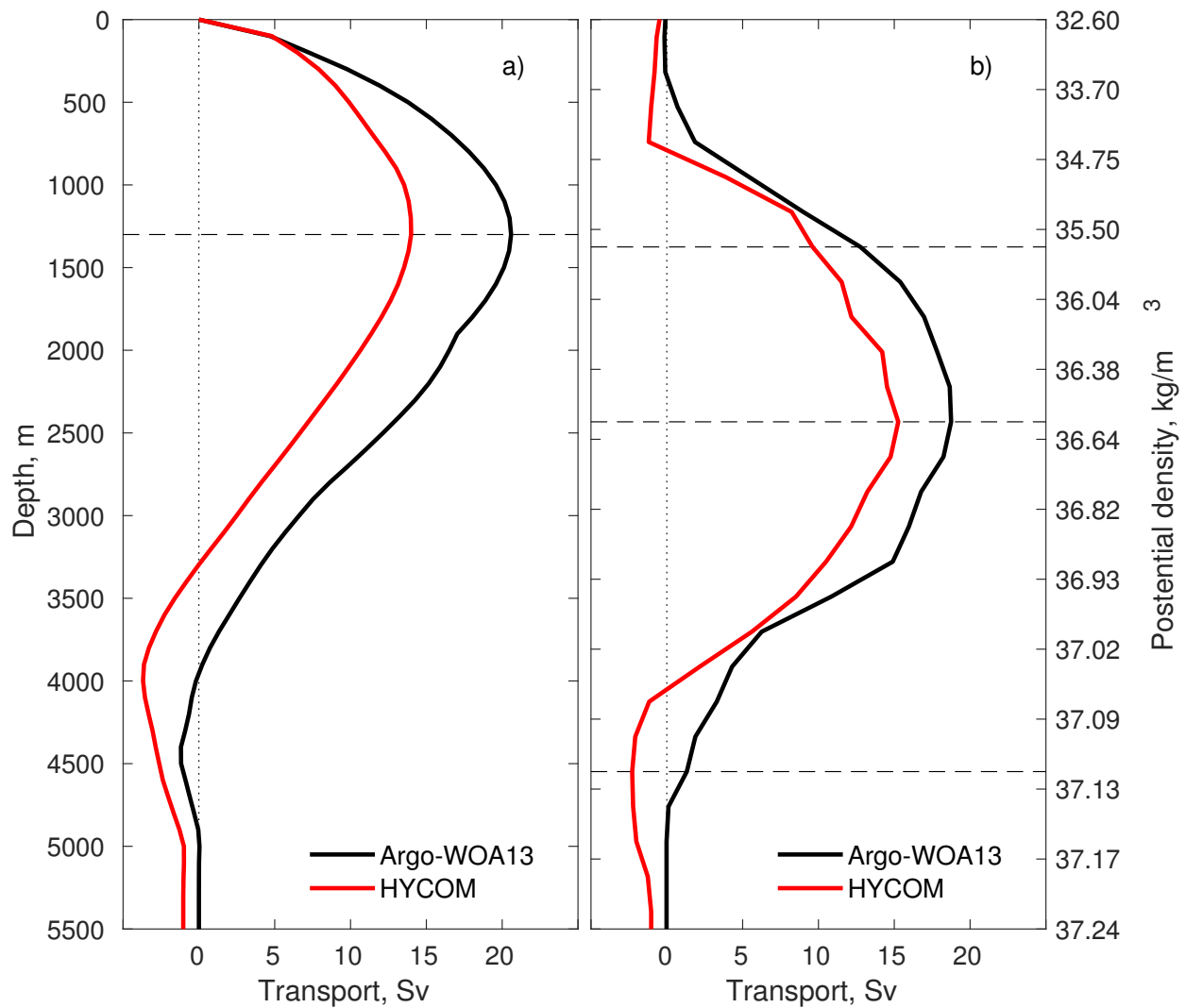


Figure 9: a) Long-term mean meridional overturning streamfunction (in Sv) at 34°S with respect to a) depth and b) potential density in σ_2 . Observations based on monthly mean Argo profile for the upper 2000 m and WOA13 below 2000 m; model results based a global 1/12° HYCOM simulation (1981-2015).

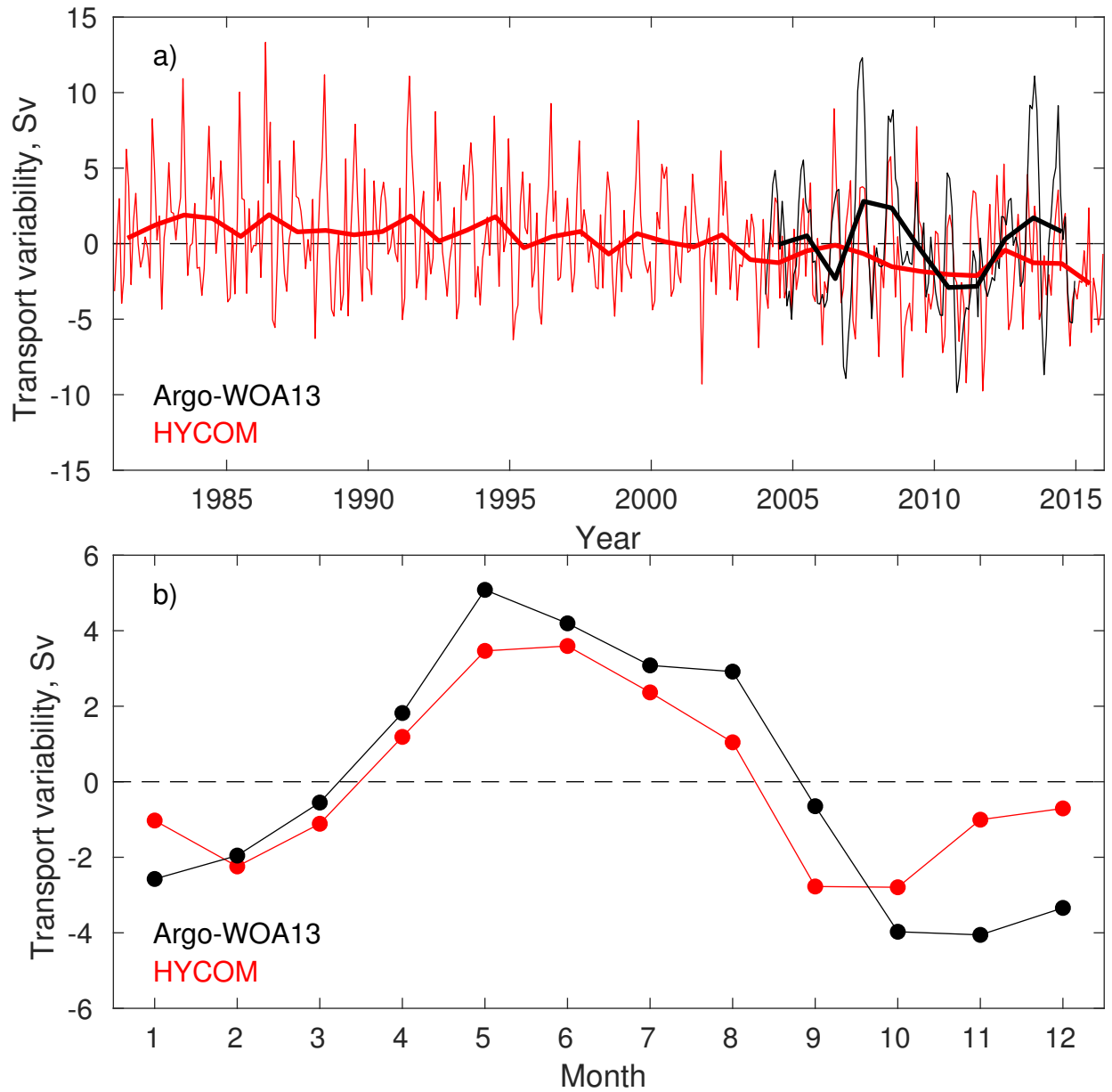


Figure 10: a) Time series of the observed (black) and modeled (red) AMOC transport at 34°S , with thin/thick lines denoting monthly/annual means, respectively; b) seasonal variability of the AMOC transports at 34°S , averaged over 2004-2014 in observations (black) and 1981-2015 in model results (red).

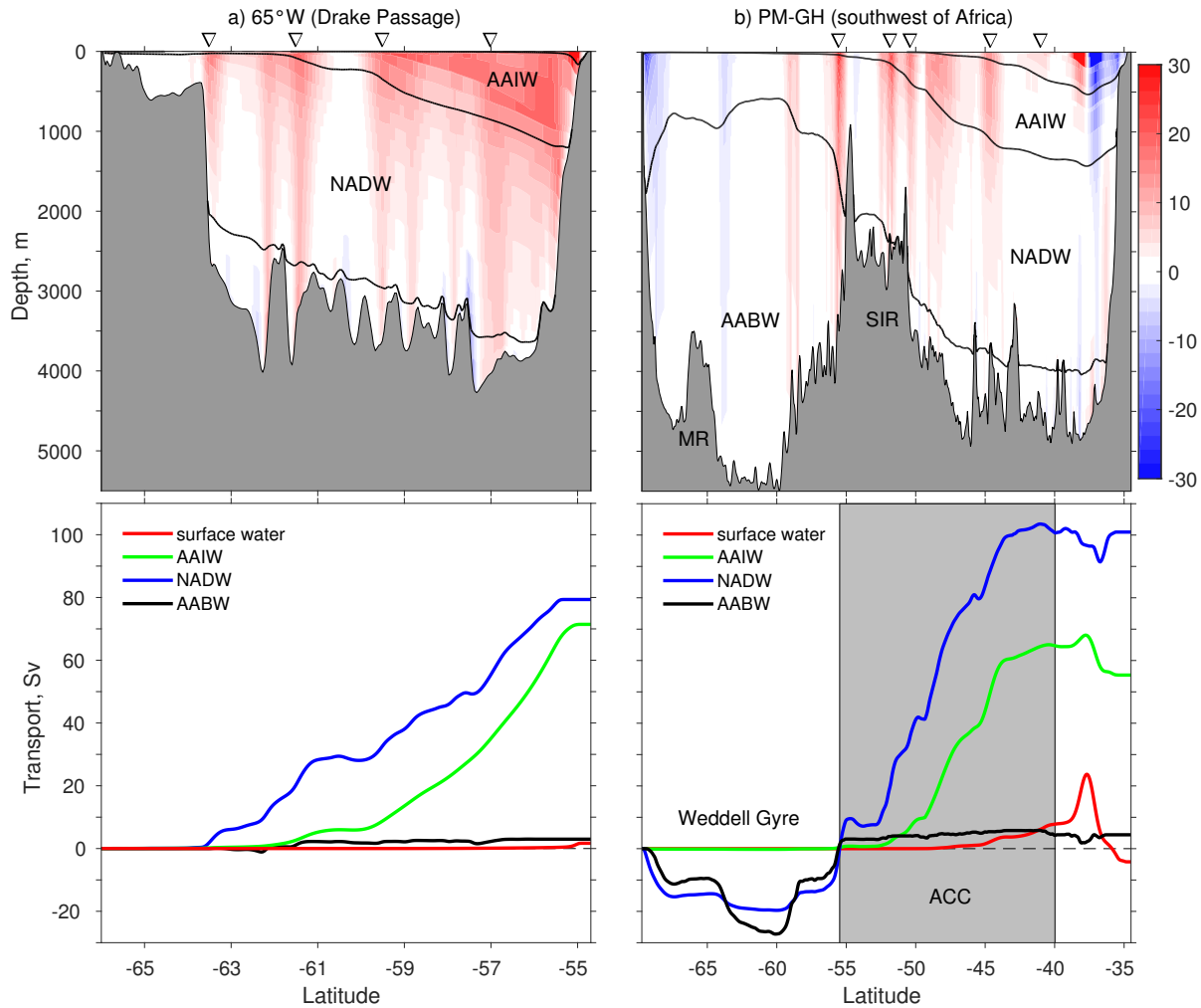


Figure 11: Modeled long-term mean zonal velocity and the corresponding four-layer volume transport in four density layers across a) 65°W in the Drake Passage and b) the Prime Meridian-Good Hope (PM-GH) transect southwest of Africa. The ∇ denote the locations of Antarctic circumpolar current (ACC) fronts, from south to north, the Southern Boundary, South ACC Front, Polar Front, Subantarctic Front, as well as the subtropical front (only in panel b). The shaded area in panel b) between 40 and 55.5°S marks the ACC regime across the PM-GH transect. Transport are accumulative northward. The four layers are near surface water ($\sigma_2 < 35.65 \text{ kg m}^{-3}$), Antarctic Intermediate Water (AAIW, $35.65 < \sigma_2 < 36.58$), North Atlantic Deep Water (NADW, $36.58 < \sigma_2 < 37.12$), and Antarctic Bottom Water (AABW, $\sigma_2 > 37.12$).

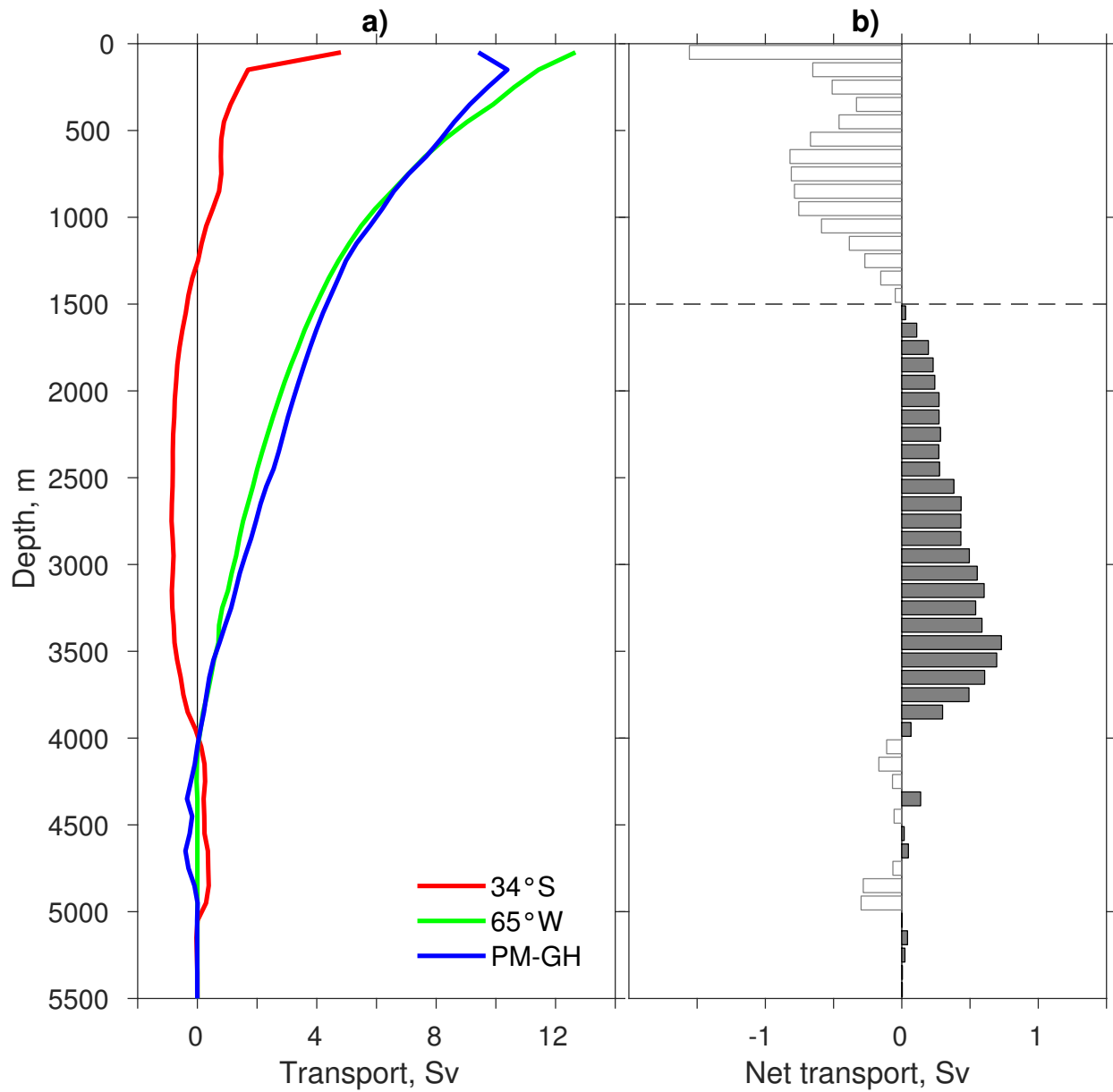


Figure 12: a) Modeled mean horizontal transports (in Sv) for very 100 m in the vertical across the 34°S, the 65°W, and the PM-GH transects. b) The net transports into the region enclosed by the three transects, with positive (negative) values indicating net transport into (out of) the region.

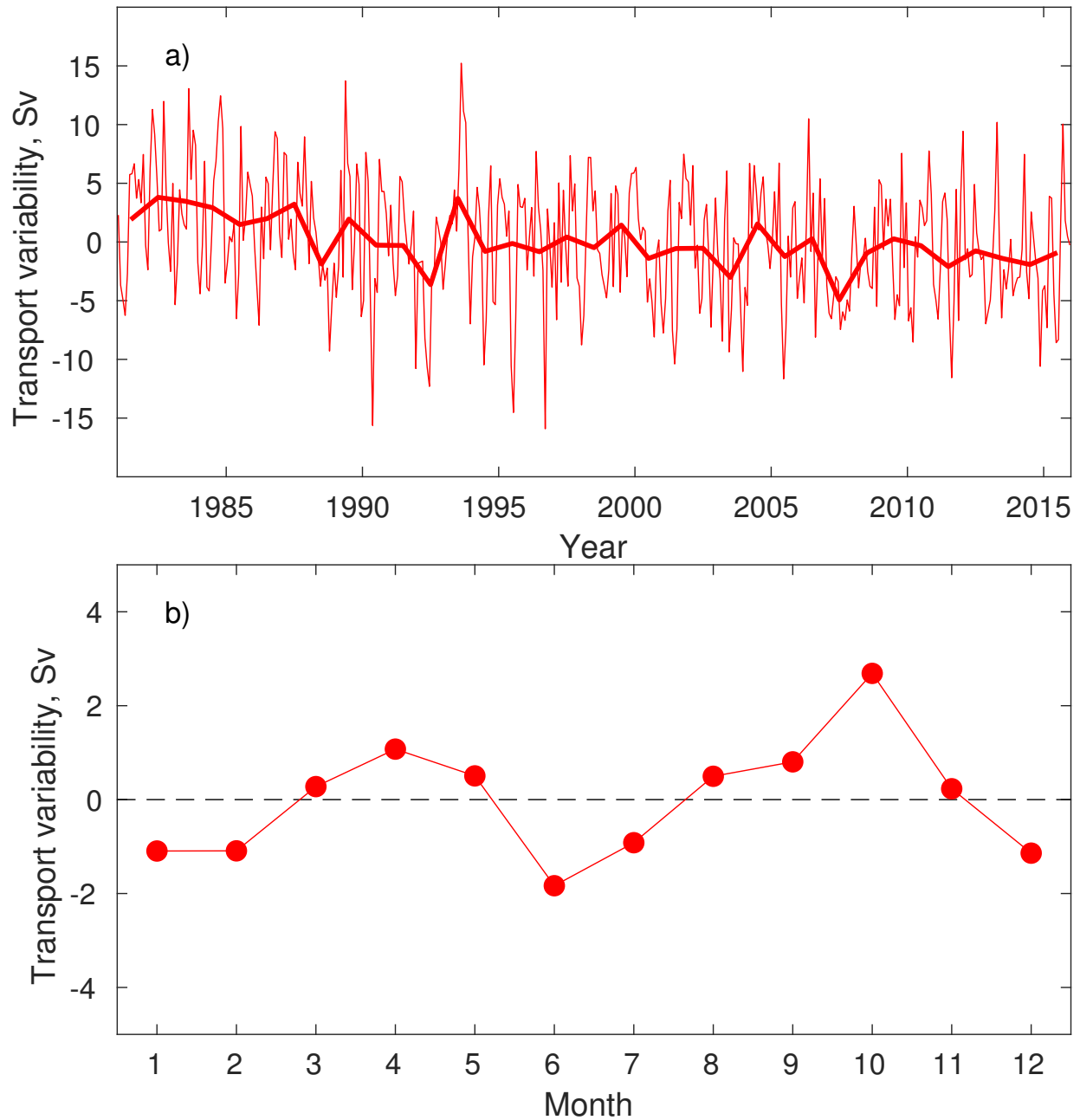


Figure 13: a) Variability of the modeled ACC transport through the Drake Passage at the 65°W , in monthly (thin) and annual (thick) means; b) Seasonal variability of the modeled ACC transports averaged over 1981-2015.

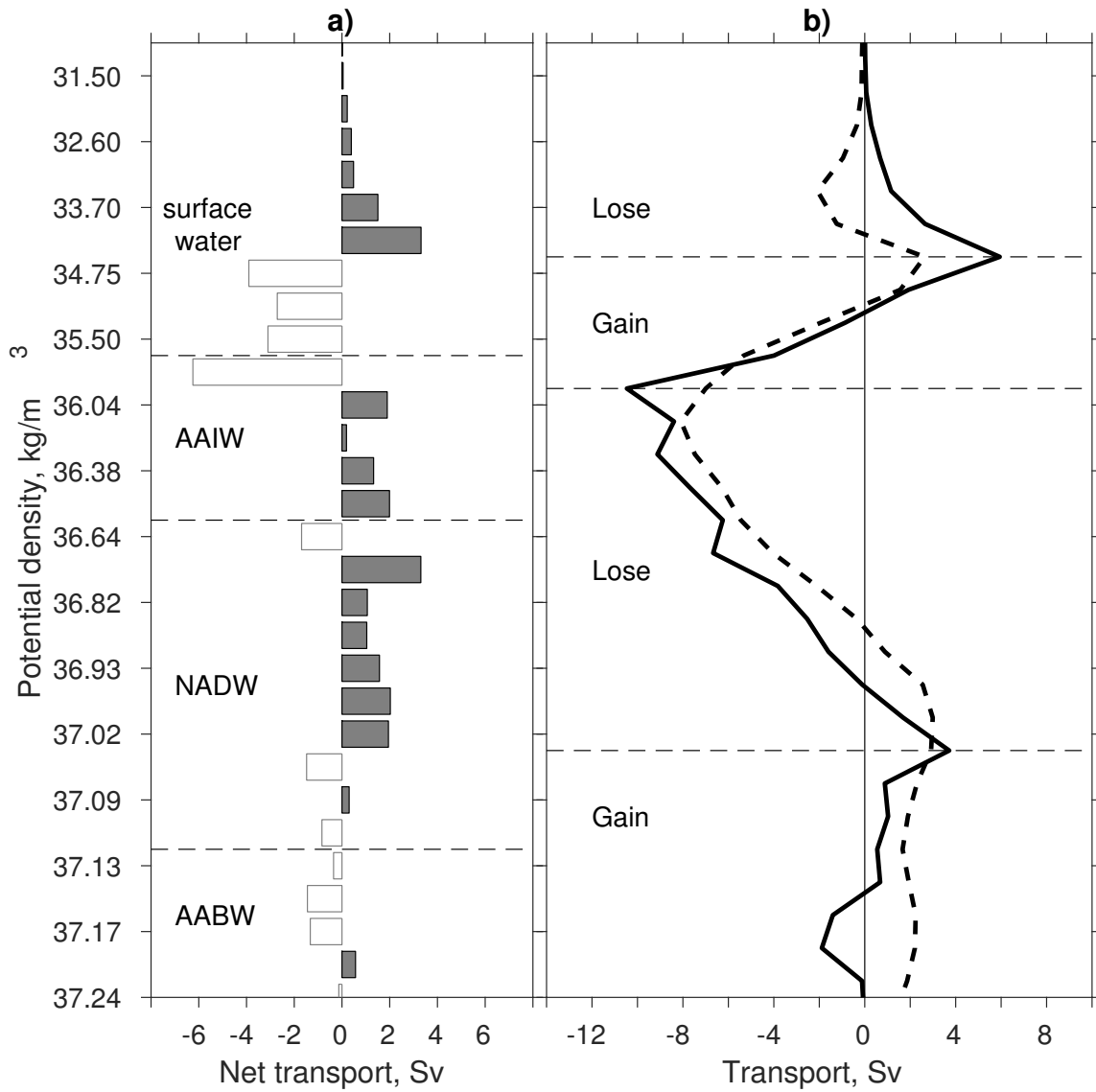


Figure 14: a) Net transports into the South Atlantic region closed by the 34°S, the 65°W, and the PM-GH transects, with respect to density layers (positive/negative values for net transport into/out of the region); b) Black line denotes the total diapycnal transformation; and dashed red line denotes the surface forced diapycnal transformation calculated from surface buoyancy fluxes.

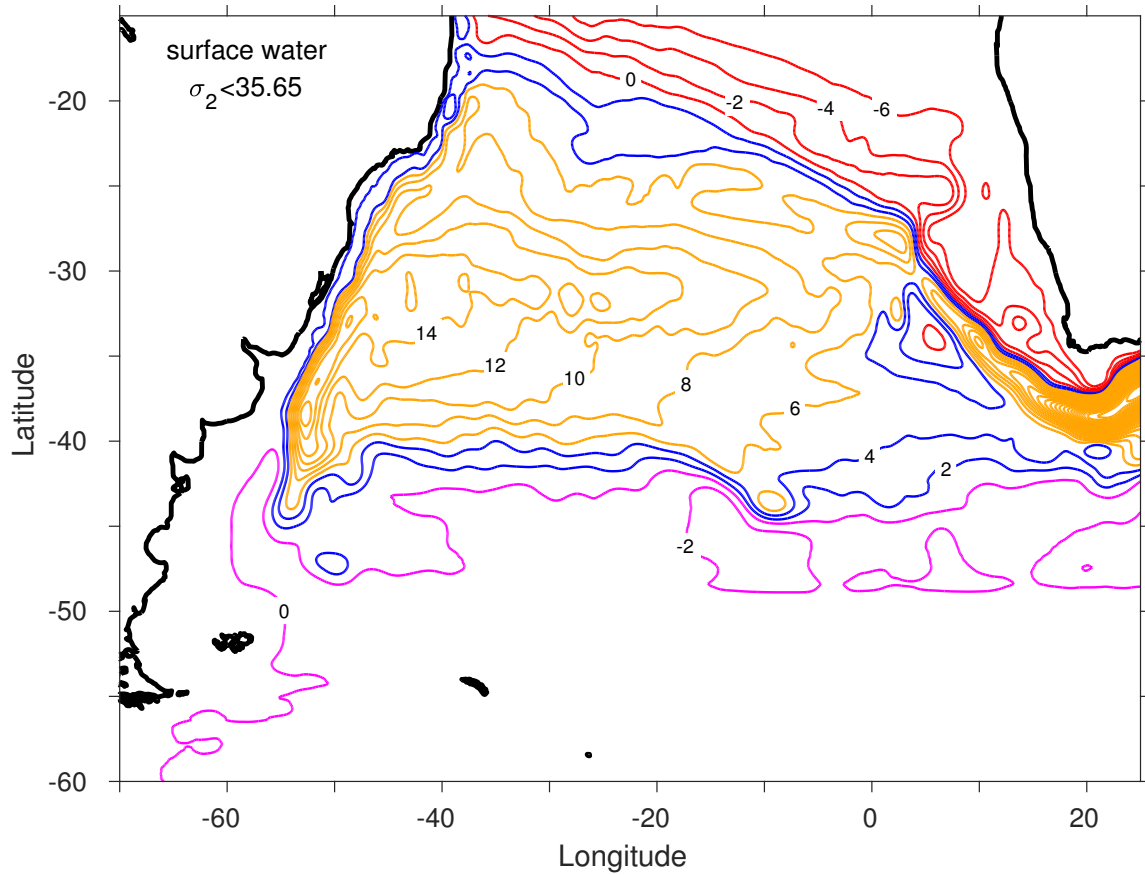


Figure 15: Modeled long-term mean horizontal transport streamfunction (in Sv) for the layer of near surface water ($\sigma < 35.65 \text{ kg m}^{-3}$). Each streamline contour is 2 Sv. Red, blue, and orange streamlines denote AMOC contribution, super gyre that flow around the South Atlantic and back to Indian Ocean, and subtropical gyre of the South Atlantic.

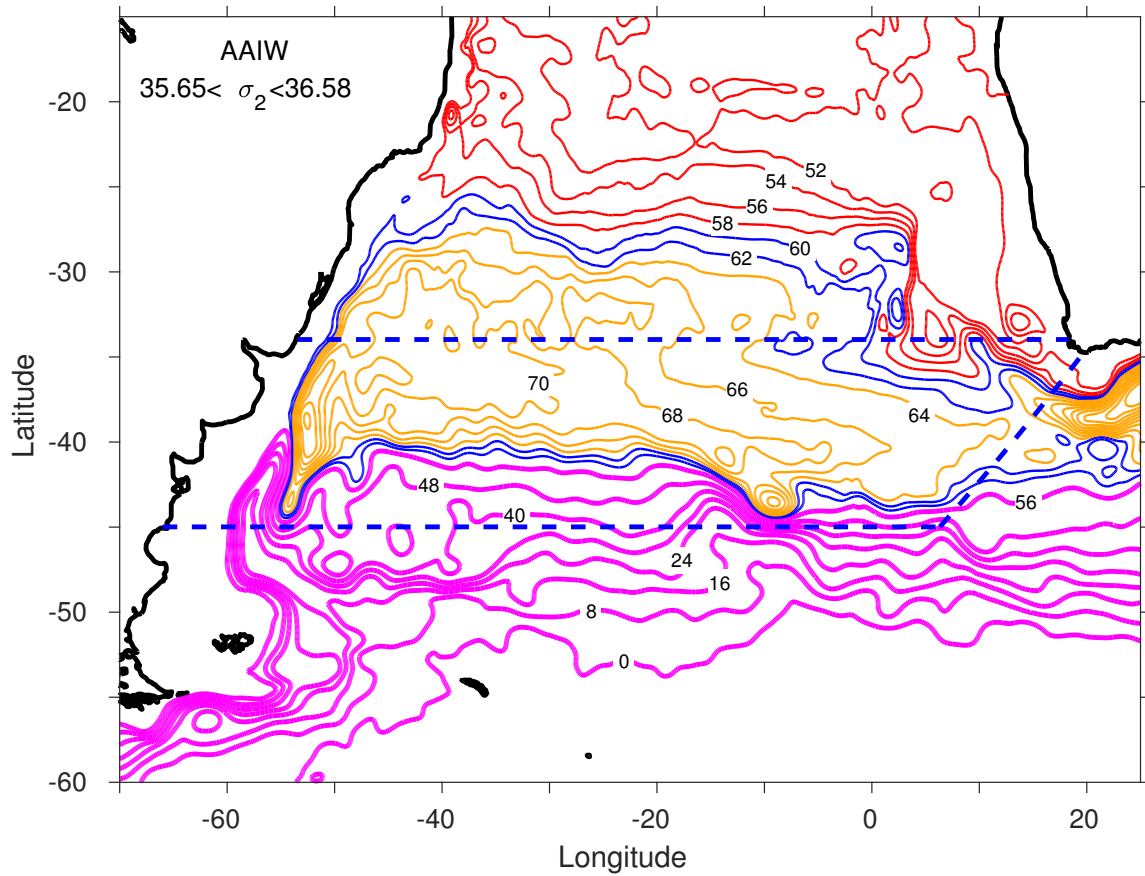


Figure 16: Modeled long-term mean horizontal transport streamfunction (Sv) for the layer of AAIW ($35.65 < \sigma < 36.58 \text{ kg m}^{-3}$). Thick pink stream lines (increment of 8 Sv) is the ACC. The Red, blue, and orange streamlines denote AMOC contribution, super gyre, and subtropical gyre of the South Atlantic (similar to Figure 15). The dashed blue lines denote 34°S , 45°S , and the GoodHope sections, across which the water properties of the northward and northwestward transports are examined Figure 17.

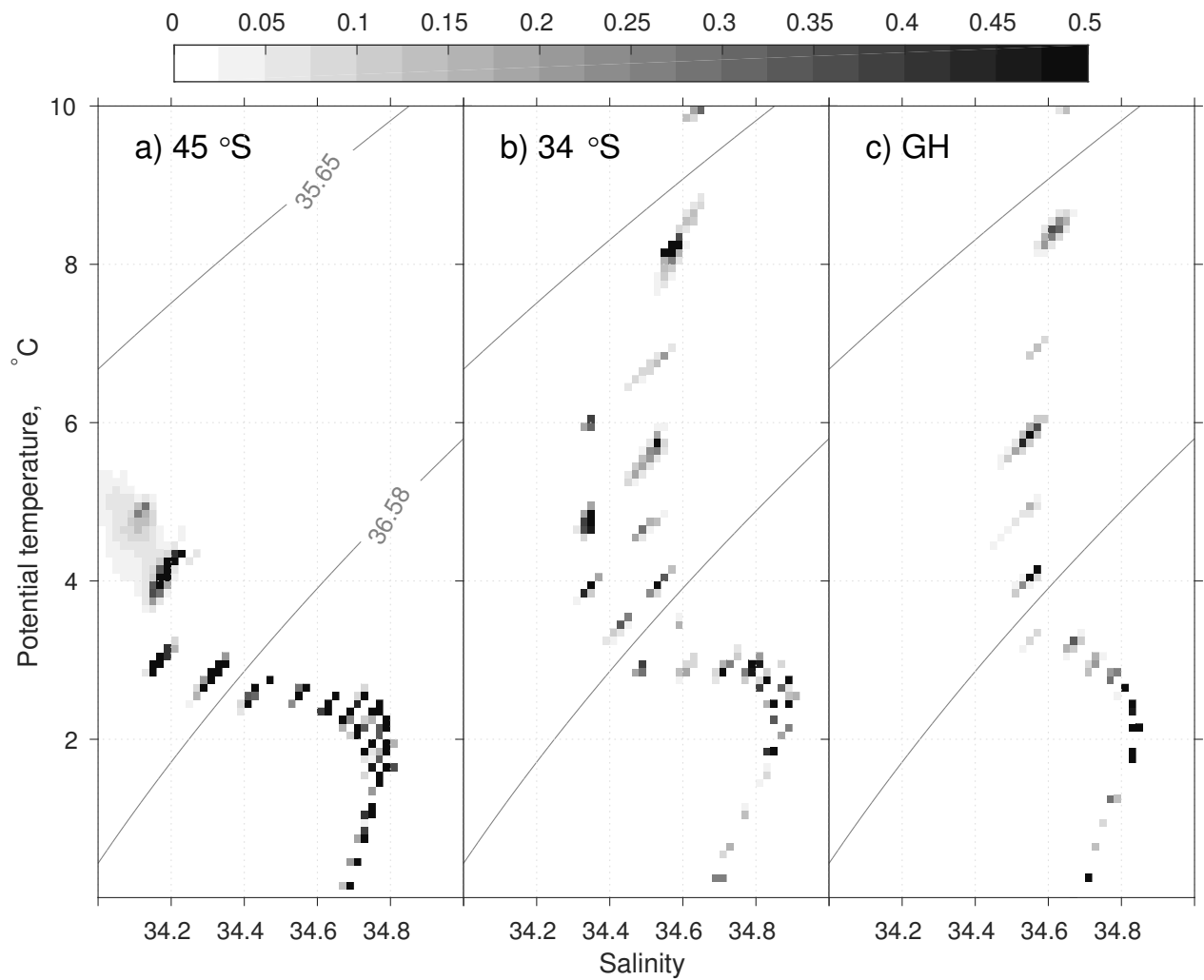


Figure 17: Modeled northward transport (in Sv) across 45°S and 34°S , and northwestward transport across the GH section, projected on potential temperature-salinity (θ -S) plane with $\Delta\theta \times \Delta S$ of $0.2^{\circ}\text{C} \times 0.04$. The isopycnal (σ_2) surfaces of 35.65 and 36.58 kg m^{-3} denote the upper and lower AAIW interfaces.

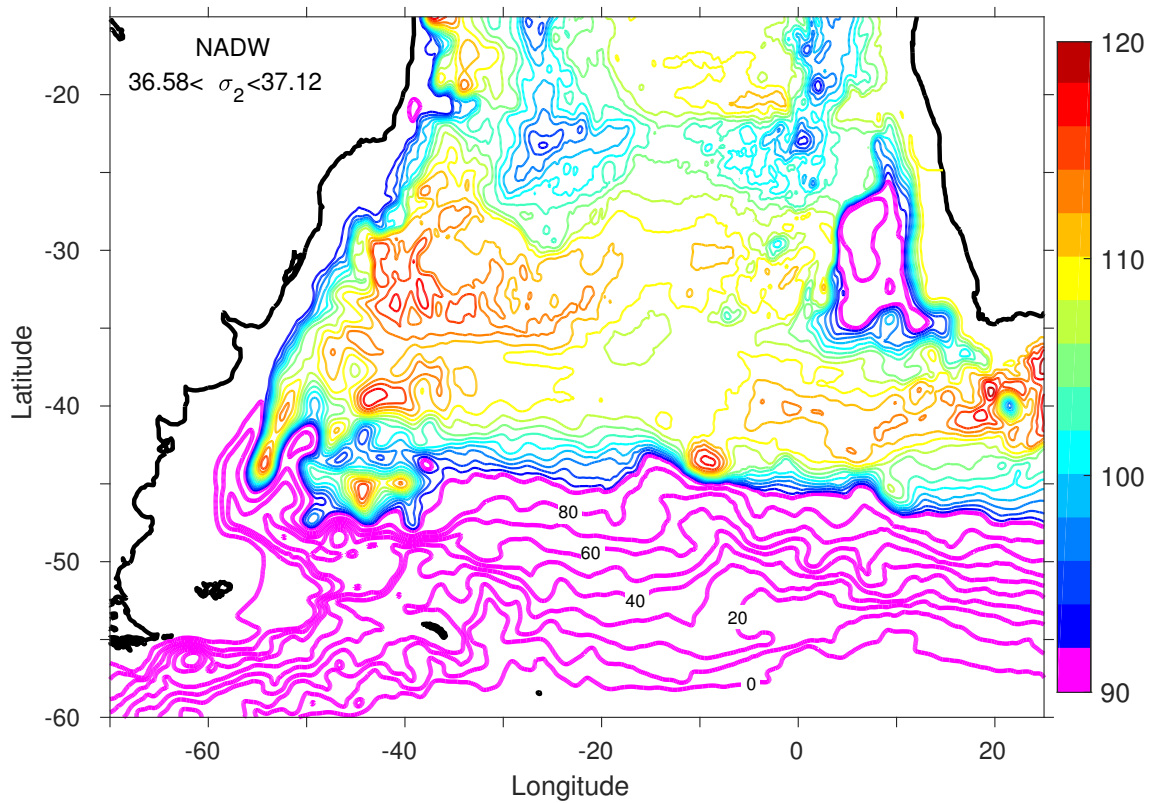


Figure 18: Modeled long-term mean horizontal transport streamfunction for the layer of NADW ($36.58 \leq \sigma_2 < 37.12 \text{ kg m}^{-3}$). Pink streamlines (10 Sv increment) indicate the eastward transport of the ACC, blue to yellow streamlines (2 Sv increment) represent the southward spreading of the NADW from north.

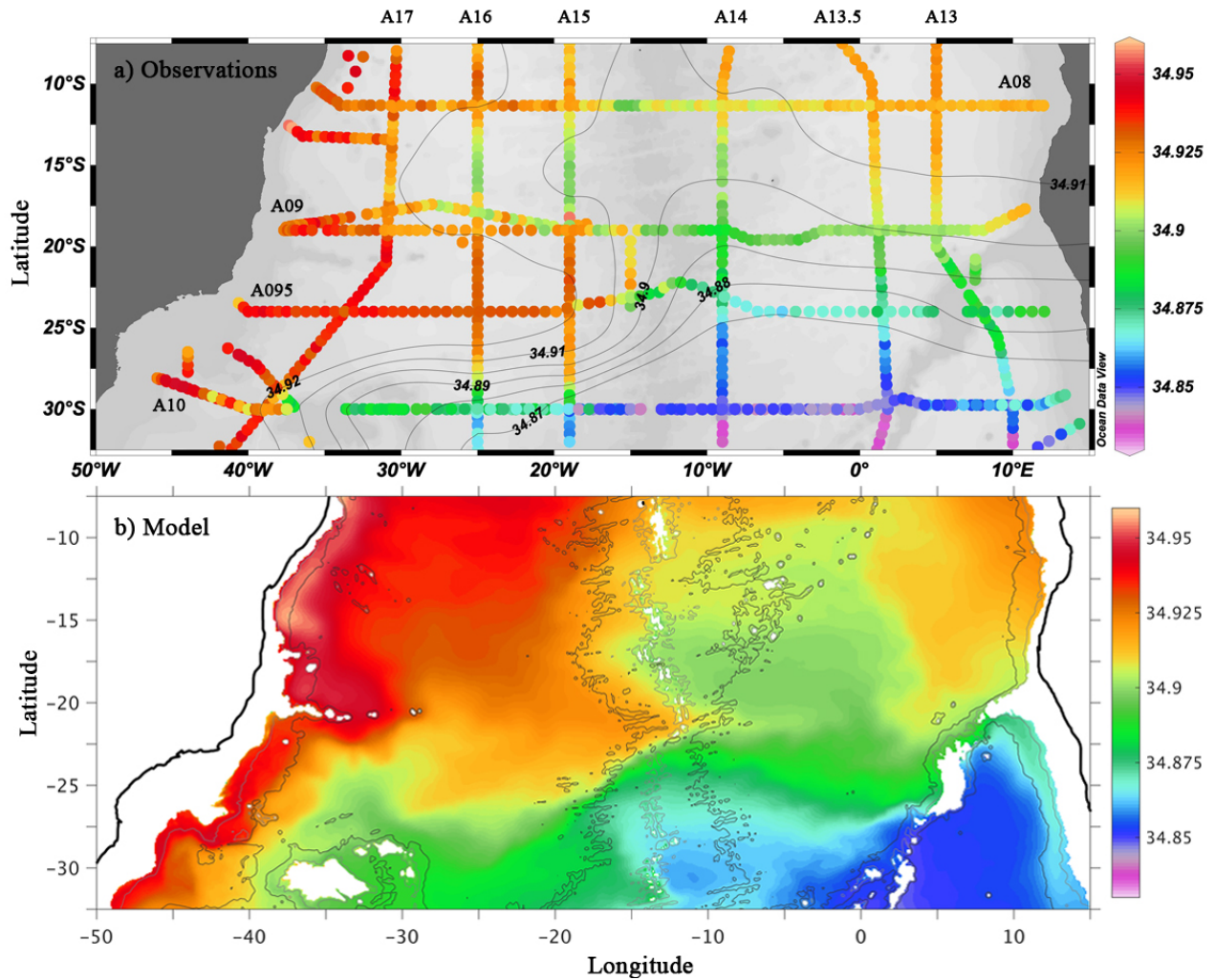


Figure 19: a) observed and model salinity distribution at 2500 m in the South Atlantic. Observation based on CTD data from GoShip program <http://www.go-ship.org>. Detailed vertical sections can be seen in the WOCE Atlas (Kiltermann et al., 2011). The results show an eastward extension of high salinity (NADW signature) between 20 and 25°S west of the mid-Atlantic Ridge (MAR), and significantly lower salinity east of MAR.

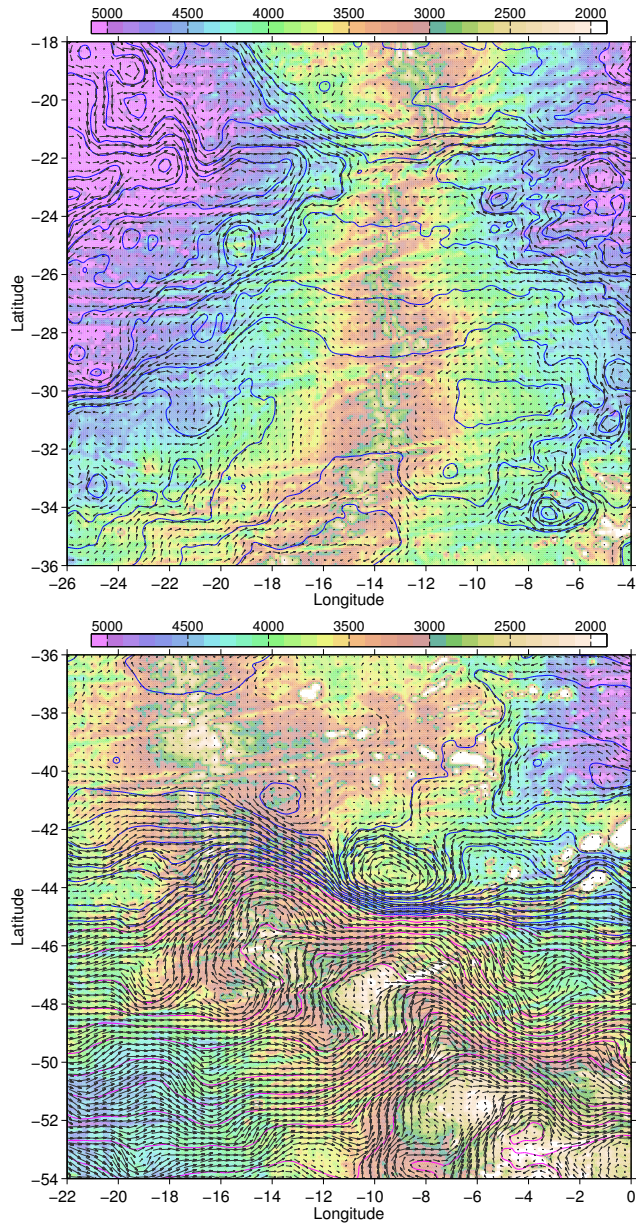


Figure 20: Zoomed view of the circulation for the density layer of NADW ($36.58 \leq \sigma_2 < 37.12 \text{ kg m}^{-3}$) across the Mid-Atlantic Ridge in the South Atlantic Ocean. The blue contours denote NADW from north and magenta streamlines denote ACC from Drake Passage.

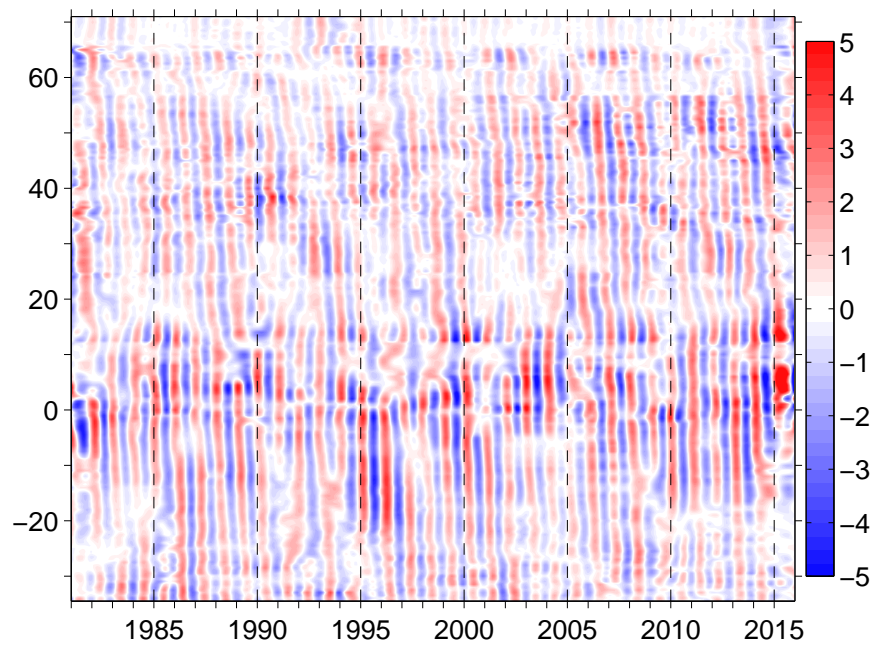


Figure 21: Modeled seasonal variability of the AMOC transports at different latitude, based on the third Intrinsic mode function (IMF) using the the ensemble empirical mode decomposition (EEMD, Huang and Wu (2008) and Wu and Huang (2009)).

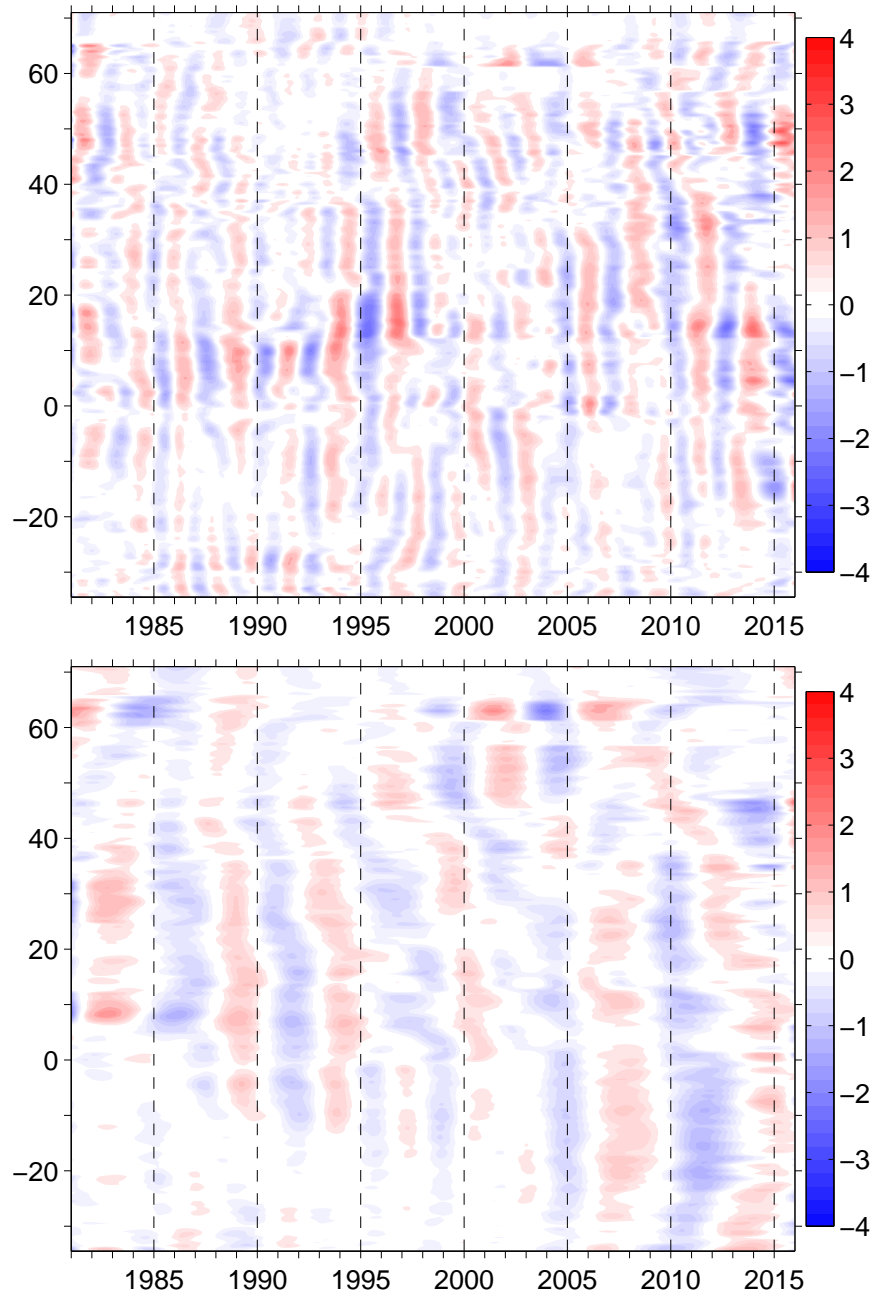


Figure 22: Modeled interannual variability of the AMOC transports at different latitude, based on the fourth and fifth Intrinsic mode function (IMF) using the ensemble empirical mode decomposition (EEMD, Huang and Wu (2008) and Wu and Huang (2009)).

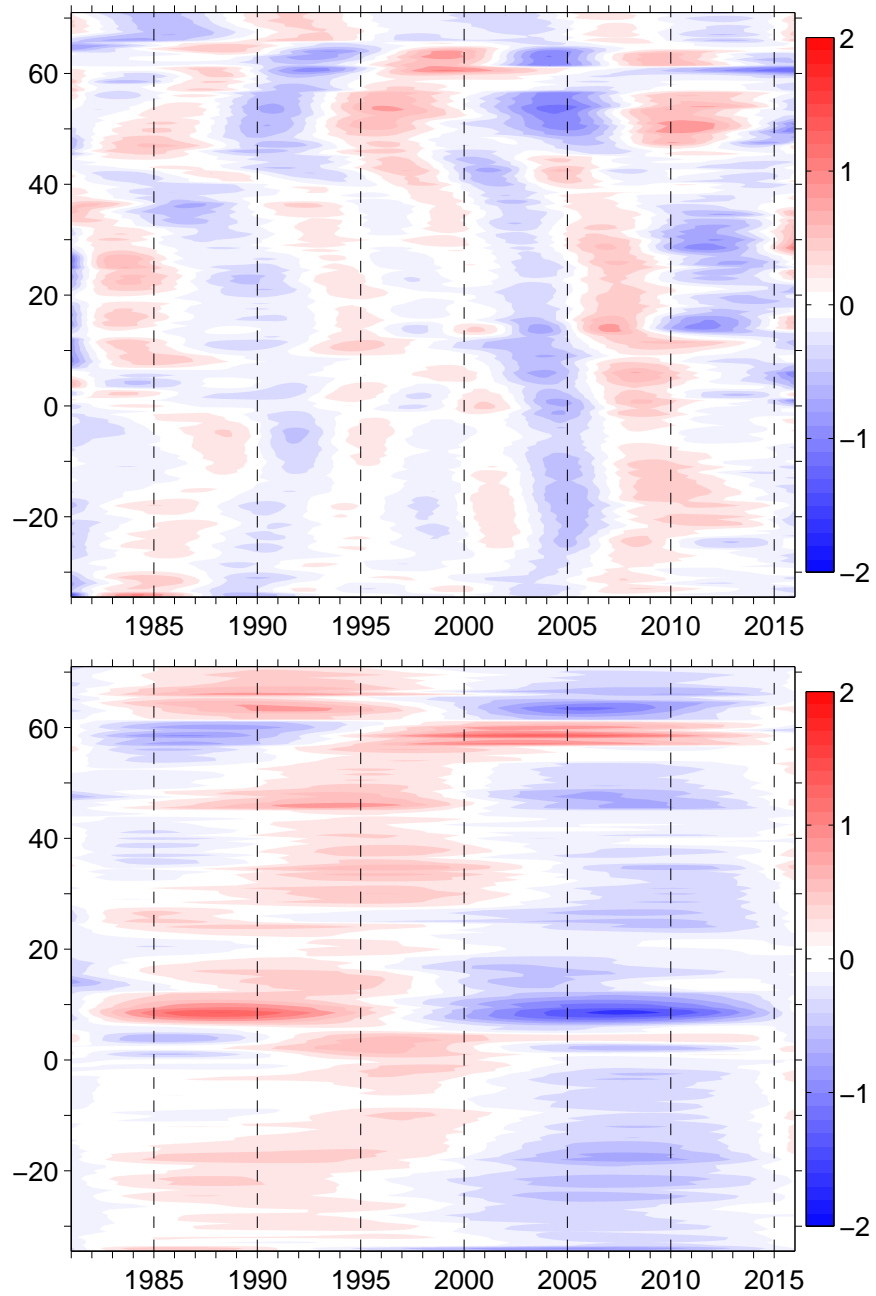


Figure 23: Modeled decadal variability of the AMOC transports at different latitude, based on the sixth and seventh Intrinsic mode function (IMF) using the ensemble empirical mode decomposition (EEMD, Huang and Wu (2008) and Wu and Huang (2009)).